

Numerical study of an individual Taylor bubble rising through stagnant liquids under laminar flow regime



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ABSTRACT

Slug flow is one of the main flow regimes encountered in multiphase flow systems especially in oil and gas production systems. In the present study, the rise of single Taylor bubble through vertical stagnant Newtonian liquid is investigated by performing complete dimensionless treatment followed by an order of magnitude analysis of the terms of equations of motion. Based on this analysis, it is concluded that Froude, Eötvös and Reynolds numbers are the sole physical parameters influencing the dimensionless slug flow equations. Using the guidelines of the order of magnitude analysis, computational fluid dynamics simulation is carried out to investigate the dynamics of Taylor bubbles in vertical pipe using the volume-of-fluid (VOF) method. Good agreement with previous experimental data and models available in the literature is established confirming that the density ratio, viscosity ratio and the initial ratio of bubble size to pipe diameter (L_{TB}/D) have minimal effect on the main hydrodynamic features of slug flow. Based on the developed results, correlations for the terminal velocity of the Taylor bubble and the dimensionless wall shear stress are proposed showing the significance of these main dimensionless parameters and support other important theoretical and experimental work available in the literature.

1. Introduction

Multiphase flows occur in a wide range of applications including natural processes, chemical processes, nuclear systems and petroleum industries. The petroleum industry is considered one of the most important applications of multiphase flow, as it could be encountered in different processes/stages such as: oil processing, oil and gas transport in pipelines, and sloshing in offshore separator devices.

For two-phase gas-liquid flow in pipes, different flow patterns can occur known as “flow pattern/flow regime”. These patterns depend on the flow rates, the geometry of the system, and inclination of the pipe (Morgado et al., 2016). Multiphase flow is classified according to the distribution of different phases building up the flow field, known as “flow regime/pattern”. Multiphase flow can be encountered in various flow patterns such as bubbly, slug, plug, annular and dispersed flow. Fluid flow investigation includes an important aspect which is the identification of the encountered flow pattern. For gas-liquid flow in pipes, one of the common and complex patterns encountered is known as “slug flow”. Slug flow is an intermitted flow between stratified and annular flow.

Flow intermittence is the main remarkable hydrodynamic characteristic causing the complex structure of slug flow which is composed of Taylor elongated bubble that occupies almost the whole cross-section of the pipe, and annular falling liquid film that might entrain many small bubbles, known as a “liquid slug”. Flooding of downstream processing facilities, severe pipe corrosion, structural instability of pipeline, and further induction of the reservoir flow oscillations, and a poor reservoir management are examples of the problems encountered as result of slugging in offshore oil and gas systems.

The prediction of the appropriate flow pattern regimes, the governing correlations, and the hydrodynamic characteristics of slug flow are essential for successful operation, simulation and optimization of any industrial applications encountering slug flow (Santos, 2007). According to the following authors, Computational fluid dynamics (CFD) has been proven to be a powerful, practical tool for the analysis and simulation of the hydrodynamic characteristics of slug flow in pipes. The main complex feature of gas-liquid slug flow is the deformable interface (Zheng and Che, 2007). The volume-of-fluid (VOF) method originally developed by Hirt and Nichols (1981) is often used to simulate complex multiphase

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Nomenclature		z^*	Dimensionless axial coordinate (–)
D	Pipe diameter (m)	<i>Greek letters</i>	
g	Acceleration due to gravity (m s^{-2})	μ	Dynamic viscosity (Pa s)
g_r	Acceleration due to gravity in radial direction (m s^{-2})	ρ	Density (kg m^{-3})
g_r^*	Dimensionless acceleration due to gravity in radial direction (–)	σ	Surface tension (N m^{-1})
g_z	Acceleration due to gravity in axial direction (m s^{-2})	σ_r	Surface tension in radial direction (N m^{-1})
g_z^*	Dimensionless acceleration due to gravity in axial direction (–)	σ_r^*	Dimensionless surface tension in radial direction (–)
L	Pipe length (m)	σ_z	Surface tension in axial direction (N m^{-1})
L_{TB}	Length of the Taylor bubble (m)	σ_z^*	Dimensionless surface tension in axial direction (–)
L_W	Length of the wake (m)	σ_θ	Surface tension in tangential direction (N m^{-1})
p	Pressure (Pa)	σ_θ^*	Dimensionless surface tension in tangential direction (–)
p^*	Dimensionless pressure (–)	δ_{LF}	Liquid film thickness (m)
r	Radial direction (m)	τ	Shear stress (Pa)
r^*	Dimensionless radial direction (–)	τ_W	Wall shear stress (Pa)
r_1 and r_2	Local principal radii of curvature at the bubble surface as indicated by Mao and Dukler (1990) (m)	ν	Kinematic viscosity ($\text{m}^2 \text{s}^{-1}$)
R	Pipe radius (m)	Γ_ρ	Density ratio, $\Gamma_\rho = \frac{\rho_L}{\rho_G}$
R_{TB}	Taylor bubble radius (m)	Γ_μ	Viscosity ratio, $\Gamma_\mu = \frac{\mu_L}{\mu_G}$
t	Time (s)	<i>Dimensionless groups</i>	
t^*	Dimensionless time (–)	Ar	Archimedes number, $Ar = \rho_L^2 g D^3 / \mu_L^2$
u	Velocity (m s^{-1})	Eo	Eötvös number, $Eo = \frac{g \rho_L D^2}{\sigma}$
U_∞	Velocity of a Taylor bubble rising through stagnant liquid (m s^{-1})	Fr_{UTB}	Froude number, $Fr_{UTB} = \frac{U_{TB}}{\sqrt{gD}}$
U_L	Mean liquid velocity (m s^{-1})	M	Morton number, $M = \frac{\Delta \rho g D^4}{\rho_L \sigma^3}$
U_{LF}	Velocity in the annular liquid film (m s^{-1})	N_f	Inverse viscosity number, $N_f = \rho_L (g D^3)^{0.5} / \mu_L$
U_{TB}	Taylor bubble velocity (m s^{-1})	$Re_{U_{TB}}$ or Re_{U_∞}	Reynolds number based on the velocity of the Taylor bubble, $Re_{U_{TB}} = \frac{\rho_L U_{TB} D}{\mu_L}$
v_r	Velocity component in radial direction (m s^{-1})	$Re_{U_{LF}}$	Reynolds number based on the velocity of the annular liquid film, $Re_{U_{LF}} = \rho_L U_{LF} \delta_{LF} / \mu_L$
v_r^*	Dimensionless velocity component in radial direction (–)	Re_{v_L}	Reynolds number based on the mean velocity of the liquid, $Re_{v_L} = \rho_L v_L D / \mu_L$
v_z	Velocity component in axial direction (m s^{-1})	<i>List of acronyms</i>	
v_z^*	Dimensionless velocity component in axial direction (–)	CFD	Computational fluid dynamics
v_θ	Velocity component in tangential direction (m s^{-1})	FRF	Fixed frame of reference
V_L	Relative liquid velocity to the bubble in moving reference frame (MRF) (m s^{-1})	MRF	Moving frame of reference
V_W	Volume of the wake (m^3)	VOF	Volume-of-fluid
x or z	Axial coordinate in 2D coordinate system (m)		
y or r	Radial coordinate in 2D coordinate system (m)		

flows including slug flow, and is powerful in tracking the interface between fluids (Fabre and Liné, 1992; Razavi and Namin, 2011; Rahimi et al., 2013; Desamala et al., 2013; Desamala et al., 2014; Fabre and Liné (1992); Razavi and Namin, 2011; Rahimi et al., 2013; Desamala et al., 2013; Desamala et al., 2014).

The hydrodynamic characteristics of gas-liquid vertical slug flow include the final shape of the Taylor bubble, Taylor bubble rises velocity, liquid film thickness, liquid film velocity, wall shear stress distribution and wake shape. Despite the conduction of extensive work in the modelling process of gas-liquid slug flow, a need for correlations based on experimental data is still required. These correlations include slug characteristics such as: Taylor bubble velocity, slug frequency, slug length, slug liquid hold up, and slug unit velocity.

In literature, since the 1940s, a significant amount of research has been done to understand the complex principles of slug flow. Starting with Dumitrescu (1943) who investigates the rise of single Taylor bubble in the stagnant liquid by applying potential flow theory and concludes that the Taylor bubble rise velocity could be given by:

$$U_{TB} = 0.351 \sqrt{gD} \quad (1)$$

Other analytical and/or experimental approaches are made later to modify the above correlation as discussed by Kang et al. (2010). A good

review on the most commonly used correlations to estimate the Taylor bubble velocity is given by Morgado et al. (2016).

One of the main complex hydrodynamic features of slug flow is the wake flow pattern. Campos and De Carvalho (1988) performs an important photographic study to investigate the wake structure of Taylor bubbles rising in stagnant liquid using different pipe diameters and liquid viscosities. They conclude that the inverse viscosity number mainly influences the wake structure and they categorise the wake flow pattern into three main groups as follows:

- Type 1: Closed axisymmetric laminar wake for: $N_f < 500$.
- Type 2: Closed asymmetric transitional wake for: $500 < N_f < 1500$.
- Type 3: Opened turbulent wake with the recirculatory flow: $N_f > 1500$.

Araújo et al. (2012) discuss the importance of other experimental studies that investigate the main complex hydrodynamic nature of slug flow. They reach number of remarkable conclusions that helped in further understanding of the problem (Polonsky et al., 1999; Van hout et al., 2002; Clanet et al., 2004; Liberzon et al., 2006; Sousa et al., 2006; Direito et al., 2017; Polonsky et al., 1999; Van hout et al., 2002; Clanet et al., 2004; Liberzon et al., 2006; Sousa et al., 2006; Direito et al., 2017).

Despite the significant effort done in most of the experimental approaches, the need for computation is essential in understanding the complex nature of slug flow problem that eliminates the experimental limitation and the difficulties while providing robust analysis and accurate results. According to the viewed literature, Computational fluid dynamics (CFD) has been proven to be a powerful, practical tool for the analysis and simulation of the hydrodynamic characteristics of slug flow in pipes.

Early attempts to numerical study of slug flow problem are made by Kawaji et al. (1997) that use the volume-of-fluid (VOF) method to numerically investigate the hydrodynamic characteristics of a Taylor bubble rising through the stagnant liquid in a vertical pipe and conclude that the bubble length does not affect the bubble terminal velocity. Based on an iterative scheme that solved the velocity and shape of Taylor bubble in a vertical tube, a different methodology is developed by Clarke and Issa (1997). They introduce a model that assume homogenous flow in the liquid slug region and thus account for the dispersed bubble in that region (Araújo et al., 2012). Against their assumption, the model shows inadequate results that they suggest that future models should use the two-fluid model for proper simulation of the dispersed gas bubbles in the liquid slug region (Ndinisa et al., 2005). Later, Bugg et al. (1998) perform a detailed study on the motion of Taylor bubbles in vertical pipes and prove that the VOF method is capable of determining the main hydrodynamic features of slug flow including the bubble shape, bubble rising velocity, liquid film thickness and average velocity in the liquid film. The results are then compared with experimental data in the literature and show good agreement. Another different approach, based on developing a new algorithm for solving the gas-liquid interface equation, is developed by Issa and Ubbink (1999).

Other significant numerical studies using the VOF method are worth to mention as those done by Ndinisa et al. (2005); Taha and Cui (2006) and Zheng and Che (2007). Lately, Lu and Prosperetti (2008) who numerically study Taylor bubble rising in the stagnant liquid by neglecting the flow in the gas using front tracking method that deals with two-phase liquids and evades the uncertain gas-liquid interface reconstruction in VOF method. They apply a procedure based on B-splines to build smooth functions best-fitting the field variables (velocities in the gas phase) over a strip straddling the gas-liquid interface along with its whole perimeter. This method show powerful tracking of the exact position and curvature of the interface that help in calculating the surface tension.

Recently, Kang et al. (2010) use a front tracking methodology to simulate the dynamics of gas slugs rising through stagnant liquids, where the finite difference method is used to discretize the governing equations. The study conclude that the density and viscosity ratios have minimal effect on the dynamics of Taylor bubbles rising in stagnant liquids, while both Eötvös number and Archimedes number have an important impact. They develop correlations for the dimensionless liquid film thickness and the dimensionless wall shear stress as a function of only Archimedes number. They also conclude that wake length and volume depend mainly on Archimedes number. Later, Araújo et al. (2013b) perform a detailed study of the dynamics of Taylor bubble rising in stagnant liquid and include a wide range of Eötvös and Morton number. They show that Kang et al. (2010) correlations are inadequate by proving that both dimensionless liquid film thickness and the dimensionless wall shear stress are function of Eötvös and Morton numbers. They also show that the wake structure is greatly influenced by both Eötvös, and Morton numbers, and not only Archimedes number as concluded by Kang et al. (2010). They are also capable of developing correlations for the wake length and volume that show good agreement with well-known experiment correlations of Campos and De Carvalho (1988).

More recently, Yan and Che (2011) investigate the hydrodynamic characteristics of single Taylor bubble rising in stagnant liquid with further consideration of the small dispersed bubbles in the liquid slug region. Their study account for the effect of small dispersed gas bubbles in liquid slug region on the flow hydrodynamics features and CO₂

corrosion rate. It is concluded that the small dispersed gas bubbles result in higher fluctuations in the liquid slug region, which subsequently increase the mass transfer and wall shear stress. Lastly, Araújo et al. (2013a) investigate the rising of two consecutive Taylor bubbles through vertical stagnant Newtonian liquids under laminar regime using the volume-of-fluid (VOF) method. The results account for bubble-bubble interaction and show the dependency of the wake on the separation distance between the bubbles. A good review on slug flow is presented by Morgado et al. (2016) that summarises all of the important correlations used in defining the problem and show the missing data that need to be further investigated.

Despite the significant numerical and theoretical data published on the rise of Taylor bubble through vertical stagnant liquid, to the authors' knowledge there is not yet a study investigating the problem using order of magnitude analysis of equations of motions. Hence, the scope of this study can be divided into two main folds:

1. Performing complete dimensionless analysis of the problem using both the Buckingham-Pi theorem and a dimensionless treatment followed by an order of magnitude analysis to the governing equations of motion in order to show the sole dimensionless parameters: $Re_{U_{TB}}$, $Fr_{U_{TB}}$, and Eo numbers and their relative merits or order of magnitudes. Based on this analysis, the main hydrodynamic features of rising of individual Taylor bubbles through stagnant Newtonian liquids are investigated by applying computational fluid dynamics (CFD) simulation using the volume-of-fluid (VOF) methodology implemented in the commercial software ANSYS Fluent.
2. Developing a correlation between $Re_{U_{TB}}$, $Fr_{U_{TB}}$, and Eo numbers based on the developed numerical results and on the guide lines of the order of magnitude analysis to predict Taylor bubble rise velocity (U_{TB}). This correlation enabled the present study to support other important theoretical and experimental work available in the literature.

2. Dimensionless equations of motion

The hydrodynamic characteristics of gas-liquid slug flow are governed by viscous, inertial, gravitational, and interfacial forces. In literature, the problem is mostly analysed into the dimensionless form using Buckingham-Pi theorem.

Morgado et al. (2016) neglect the effect of the expansion of gas during its rise and show that the dimensionless analysis of the problem results in the following form:

$$\frac{U_{TB}^2 \rho_L}{g D \Delta \rho} = f \left[Eo = \frac{g \rho_L D^2}{\sigma}, M = \frac{\Delta \rho g \mu_L^4}{\rho_L^2 \sigma^3}, \Gamma_\mu = \frac{\mu_L}{\mu_G}, \Gamma_\rho = \frac{\rho_L}{\rho_G}, \frac{L_{TB}}{D} \right] \quad (2)$$

where the LHS of relation (2) represents the ratio between the inertia and gravitational forces and known by Froude number ($Fr_{U_{TB}}$). Relation (2) can be modified by introducing Reynolds number based on bubble velocity ($Re_{U_{TB}}$) which is the ratio between the inertial forces and viscous forces. Eötvös number (Eo) is the ratio between gravitational forces, and surface tension forces, and Morton number (M) is known by property group as it only contains the properties of the fluid (Araújo et al., 2012).

It can be shown in this respect that Morton number (M) does not represent any peculiar physical quantity as it can be written as:

$$M = \left(\frac{Fr_{U_{TB}}}{Re_{U_{TB}}} \right)^4 Eo^3 = \frac{Eo^3}{N_f^4} \quad (3)$$

where the inverse viscosity number (N_f) is defined as:

$$N_f = \frac{Re_{U_{TB}}}{Fr_{U_{TB}}} = \frac{\rho_L (g D^3)^{0.5}}{\mu_L} \quad (4)$$

In fact, the inverse viscosity number (N_f) can be interpreted physically as the ratio between gravity force and the viscous force. It is worth noting in this respect that N_f and M cannot generally replace $Re_{U_{TB}}$ to

judge whether the flow is in its laminar or turbulent regime.

It can be easily shown here that using the Buckingham-Pi theorem can lead to the same form of the dimensionless groups with Morton number replaced by Reynolds number that is given by:

$$Fr_{U_{TB}} = \frac{U_{TB}}{\sqrt{gD}} = f \left[Eo = \frac{g\rho_L D^2}{\sigma}, Re_{U_{TB}} = \frac{\rho_L U_{TB} D}{\mu_L}, \Gamma_\mu = \frac{\mu_L}{\mu_G}, \Gamma_\rho = \frac{\rho_L}{\rho_G}, \frac{L_{TB}}{D} \right] \quad (5)$$

The main governing equations of the problem are the continuity and momentum equations. Dimensionless analysis of these equations followed by the order of magnitude analysis in all directions and on the boundaries is given in details in Appendix 1. The main results are: the radial velocity component (v_r^*) should be of order $\left(\frac{D}{L_{TB}}\right)$ in order to keep the continuity equation intact without any approximation; the pressure gradient in the radial direction should be of order $\left(\frac{D}{L_{TB}}\right)^2$; for all terms in the momentum equations and on the boundaries to remain of the same order of magnitude, $Fr_{U_{TB}}$ should be of order (1), $Re_{U_{TB}}$ and Eo both should be of order $\left(\frac{L_{TB}}{D}\right)$. Added to this, referring to Appendix 1, it has been shown that the inverse viscosity number is in fact a modified Reynolds number provided that the characteristic velocity for stagnant fluid given by: $V_s = (gD)^{1/2}$.

In conclusion, the dimensionless analysis show that any analysis of the problem should include the dimensionless groups: $Fr_{U_{TB}}$, Eo , and $Re_{U_{TB}}$ or N_f . In addition, it is worth to mention that there is no need to using a dimensionless number like Morton number which does not seem to have a direct physical meaning.

Thus, the problem in the present study is analysed in terms of six main dimensionless groups, namely: Eötvös number (Eo), Froude number ($Fr_{U_{TB}}$), Reynolds number based on bubble velocity ($Re_{U_{TB}}$), density ratio

(Γ_ρ), viscosity ratio (Γ_μ) and the initial ratio of bubble size to pipe diameter (L_{TB}/D). The present study examines the effects of density ratio (Γ_ρ), viscosity ratio (Γ_μ) and the initial ratio of bubble size to pipe diameter (L_{TB}/D) for the sake of supporting other previous experimental and numerical works in the literature. This is to allow dealing carefully with the main influencing parameters: Eötvös number (Eo) and Reynolds number ($Re_{U_{TB}}$).

3. CFD model development

The flow domain is constructed and solved using the volume-of-fluid (VOF) methodology implemented in the computational fluid dynamics software package, ANSYS Fluent (Release 15.0). In all simulated cases, a uniform grid of quadrilateral control elements is applied. Different grids depending on domain dimensions have been tested to check solution convergence. The present simulation has been performed for unsteady flow with constant fluid properties. The two phases are assumed as incompressible, viscous, immiscible, and not penetrating each other. The flow regime could be laminar which means the transition or the turbulent is depending on the value of Reynolds number in different flow regions of the slug flow domain.

Fig. 1 shows that the flow regions in vertical slug flow that can be divided into three main regions according to the definition of Reynolds number, into flow in main liquid region (liquid slug), Re_{U_∞} or $Re_{U_{TB}}$, flow in liquid film (annular film), $Re_{U_{LF}}$, and flow near bubble wake, Re_{V_L} . The definition of Reynolds number according to each region is given by:

$$Re_{U_{TB}} = \rho_L U_{TB} D / \mu_L \quad (6)$$

$$Re_{U_{LF}} = \rho_L U_{LF} \delta_{LF} / \mu_L \quad (7)$$

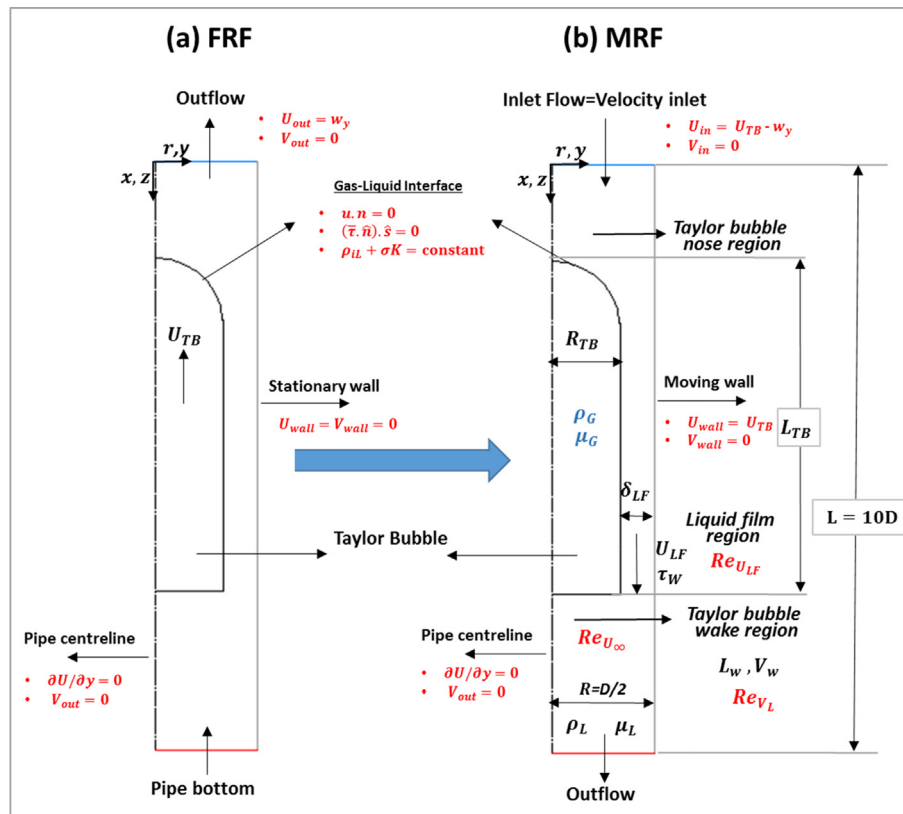


Fig. 1. Schematic representation of the computational domain and the main hydrodynamic features of a single Taylor bubble rising through stagnant liquid showing the different flow regions based on the definition of Reynolds number (Mayor et al., 2007), the initial and boundary conditions, the change from a fixed reference frame system (a) to a moving reference frame system (b).

$$Re_{V_L} = \rho_L V_L D / \mu_L \tag{8}$$

where: U_{TB} is the velocity of Taylor bubble rising through stagnant liquid according to the definition of (Nicklin et al., 1962), U_{LF} is the absolute velocity of the liquid in the stabilized liquid film of thickness δ_{LF} , and V_L is the relative liquid velocity to the bubble in moving reference frame (MRF).

For the case of Taylor bubble rising in stagnant liquid only two parameters, namely: $Re_{U_{TB}}$, and $Re_{U_{LF}}$, are used to characterize flow type into: laminar, transient, or turbulent flow regime. Various experimental work is done to investigate the characteristic Reynolds number for the onset of transition in the flow regime in each region discussed above. For instance, to ensure laminar flow regime in the main liquid region, $Re_{U_{TB}}$ should be less than 2100 as reported by (Fulford, 1964; Fréchou, 1986; Mayor et al., 2007).

Also, the range of Reynolds numbers for the transitional region in the liquid film region is $[250 : 400] < Re_{U_{LF}} < 800$ as indicated by Fulford (1964). It should be pointed out that the transition criterion from laminar into turbulent flow is not clear enough and need to be further

investigated. For instance, the wake pattern is identified into laminar or turbulent flow according to the value of inverse viscosity number (N_f). Nevertheless, according to the problem definition it should be mainly, in terms of Reynolds number. However according to data in literature, to ensure laminar flow regime in the main liquid region for the present study, $Re_{U_{TB}}$ ought to be less than 200. Regarding the flow in the liquid film, $Re_{U_{LF}}$ never exceeded 30 which ensures that the developed liquid films are entirely under laminar regime.

In the present model, the fluids share a well-defined interface, and hence the volume-of-fluid (VOF) method for two-phase flow has been selected in CFD software ANSYS Fluent (Release 15.0). The VOF model developed by Hirt and Nichols (1981) is a surface-tracking technique applied to a fixed Eulerian mesh. This model is designed for two or more immiscible fluids to track the interface between them. This model solves a single set of momentum equation that is shared by the two fluids, and the volume fraction of each of the fluids in each computational cell is followed throughout the domain. The finite volume method is used to discretize the governing equations. Details of the governing equations

Table 1
Simulation cases and their corresponding results.

Input Parameters					Predicted Values							
Case Number	E_o	$\Gamma_\rho = \frac{\rho_L}{\rho_G}$	$\Gamma_\mu = \frac{\mu_L}{\mu_G}$	$\left(\frac{L_{TB}}{D}\right)_i$	$Fr_{U_{TB}}$	$Re_{U_{TB}}$	$\left(\frac{L_{TB}}{D}\right)_p$	Flow in liquid film region			Flow in wake region	
								$\frac{\delta_{LF}}{D}$	U_{LFmax}	$\left(\frac{\tau_w}{\rho_L g D}\right)_{max}$	L_w/D	V_w/D^3
Effect of density ratio (Γ_ρ)												
1	66	6E+01	6.67E+03	3	0.290	24.3	3.1	0.1250	0.1870	0.1074	0	0
2	66	9.98E+02	6.67E+03	3	0.290	24.3	3.1	0.1243	0.1086	0.109	0	0
3	66	2E+02	6.67E+03	3	0.287	24.1	3.1	0.1246	0.1873	0.1084	0	0
Effect of density ratio (Γ_μ)												
4	66	9.98E+02	6.67E+01	3	0.288	24.2	3.12	0.1251	0.1848	0.1071	0	0
5	66	9.98E+02	6.67E+02	3	0.289	24.3	3.1	0.1244	0.1863	0.1085	0	0
6	66	9.98E+02	6.67E+03	3	0.290	24.3	3.1	0.1243	0.1086	0.109	0	0
Effect of L_{TB}/D												
7	66	9.98E+02	6.67E+03	2	0.2898	24.3	2.14	0.1246	0.1083	0.1083	0	0
8	66	9.98E+02	6.67E+03	3	0.2898	24.3	3.1	0.1243	0.1086	0.109	0	0
9	66	9.98E+02	6.67E+03	4	0.2898	24.3	4.1	0.1244	0.1087	0.1087	0	0
10	66	9.98E+02	6.67E+03	10	0.2847	23.9	9.9	0.1232	0.1084	0.1084	0	0
Effect of Reynolds number ($Re_{U_{TB}}$)												
11	66	9.98E+02	6.67E+03	3	0.246	12	3.04	0.138	0.1437	0.1188	0	0
12	66	9.98E+02	6.67E+03	3	0.292	29	3.09	0.1183	0.1993	0.2350	0	0
13	66	9.98E+02	6.67E+03	3	0.307	46	3.08	0.1063	0.1041	0.0948	0.3775	0.06779
14	66	9.98E+02	6.67E+03	3	0.315	63	3.11	0.0982	0.2596	0.0882	0.5717	0.1305
15	66	9.98E+02	6.67E+03	3	0.320	80	3.1	0.0922	0.2815	0.0831	0.6542	0.1786
16	66	9.98E+02	6.67E+03	3	0.324	97	3.12	0.0875	0.2995	0.0791	0.7341	0.2212
17	66	9.98E+02	6.67E+03	3	0.325	114	3.14	0.0822	0.3185	0.0788	0.8087	0.2571
18	66	9.98E+02	6.67E+03	3	0.328	131	3.23	0.0799	0.3281	0.0728	0.8824	0.2870
19	66	9.98E+02	6.67E+03	3	0.329	148	3.21	0.0779	0.3413	0.0708	0.9023	0.3048
20	66	9.98E+02	6.67E+03	3	0.330	165	3.27	0.0746	0.3914	0.0689	0.9331	0.32780
Effect of Eötvös number (E_o)												
21	6	9.98E+02	6.67E+03	3	0.031	2.6	3.06	0.0440	0.0540	0.1290	0	0
22	10	9.98E+02	6.67E+03	3	0.118	9.9	3.14	0.0765	0.1179	0.1510	0	0
23	20	9.98E+02	6.67E+03	3	0.216	18.2	3.1	0.1033	0.1631	0.1360	0	0
24	40	9.98E+02	6.67E+03	3	0.272	22.9	3.06	0.1221	0.1797	0.1069	0	0
25	66	9.98E+02	6.67E+03	3	0.290	24.3	3.1	0.1243	0.1864	0.1086	0	0
26	70	9.98E+02	6.67E+03	3	0.290	24.4	3.1	0.1245	0.1873	0.1087	0	0
27	80	9.98E+02	6.67E+03	3	0.292	24.5	3.1	0.1247	0.1884	0.1089	0	0
28	100	9.98E+02	6.67E+03	3	0.295	24.8	3.12	0.1251	0.1907	0.1092	0.5511	0.00561
29	120	9.98E+02	6.67E+03	3	0.296	24.9	3.12	0.1253	0.1916	0.1093	0.1796	0.01633
30	140	9.98E+02	6.67E+03	3	0.297	25.0	3.14	0.1253	0.1921	0.1094	0.2323	0.01947
31	150	9.98E+02	6.67E+03	3	0.298	25.0	3.14	0.1254	0.1924	0.1095	0.2227	0.02507
32	160	9.98E+02	6.67E+03	3	0.298	25.1	3.14	0.1254	0.1931	0.1097	0.2290	0.02702
33	170	9.98E+02	6.67E+03	3	0.298	25.1	3.14	0.1255	0.1930	0.1096	0.2552	0.0246
34	180	9.98E+02	6.67E+03	3	0.298	25.1	3.14	0.1251	0.1940	0.1102	0.2556	0.0318
35	200	9.98E+02	6.67E+03	3	0.299	25.1	3.16	0.1251	0.1941	0.1101	0.2730	0.03241
36	250	9.98E+02	6.67E+03	3	0.301	25.3	2.98	0.1249	0.1951	0.1109	0.3734	0.10823
37	300	9.98E+02	6.67E+03	3	0.301	25.3	3	0.1250	0.1955	0.1110	0.4246	0.08719
38	350	9.98E+02	6.67E+03	3	0.299	25.1	2.9	0.1247	0.1944	0.1107	0.4875	0.17962
39	400	9.98E+02	6.67E+03	3	0.299	25.2	2.89	0.1248	0.1926	0.1096	0.5867	0.13278
40	500	9.98E+02	6.67E+03	3	0.301	25.3	3.04	0.1253	0.1918	0.1094	0.5921	0.11758
41	600	9.98E+02	6.67E+03	3	0.297	24.9	2.86	0.1248	0.1953	0.1109	0.3346	0.07387
42	700	9.98E+02	6.67E+03	3	0.297	24.9	2.88	0.1247	0.1945	0.1107	0.4989	0.09641

and the treatment of the interface can be obtained from [Fluent \(2015\)](#). The continuum surface force (CSF) of [Brackbill et al. \(1992\)](#) is used to account for the surface tension effects.

3.1. Model geometry and boundary conditions

The solution domain is a vertical pipe with diameter (D) and length (L) with symmetry along the centreline of the pipe. In order to minimise computational time and effort, all the simulations were performed in axisymmetric flow situation, assuming symmetry about the centerline of the pipe. This assumption is adequate and based on the laminar state of flow. In all simulation cases, the length of domain is eleven times larger than pipe diameter to avoid disturbance of the continuous phase (liquid phase) at the entrance and the exit regions, and to ensure that a uniform velocity profile is restored. [Fig. 1](#) shows a schematic representation of the computational domain for the present problem. The initial bubble shape is a cylinder connected to a hemisphere with the same radius giving an overall bullet shape of Taylor bubble. The length and radius of the Taylor bubble are given by: L_{TB} , and R_{TB} respectively. The initial ratio of Taylor bubble length to pipe diameter (L_{TB}/D)_i is an input parameter prior to simulation and it is fixed to value of 3 in most of the investigated cases except cases 7, 9, and 10 in [Table 1](#). This initial shape is simulated until a steady bubble shape is reached. Different bubble shapes are tested and final steady shape of bubble is found to be similar but this only affects the solution convergence.

The simulation is performed by attaching a reference frame to the rising Taylor bubble. Enabling moving reference frame (MRF) in the simulation, causes the rising Taylor bubble to be stationary and the pipe wall moves downwards with a velocity equal to that of the bubble ([Mao and Dukler, 1990](#)). The transformation of the boundary conditions using MRF is given in [Fig. 1](#). The initial guess of Taylor bubble velocity (U_{TB}) is estimated according to the general correlation of [Wallis \(1969\)](#), which is given by:

$$Fr = \frac{U_{TB}}{\sqrt{gD}} = 0.345 \left(1 - e^{-\frac{0.01N_f}{0.345}} \right) \left(1 - e^{-\frac{3.37-Eo}{m}} \right)$$

where:

$$m = \begin{cases} 25, & N_f < 18 \\ 69N_f^{-0.35}, & 18 < N_f < 250 \\ 10, & N_f > 250 \end{cases} \quad (9)$$

Once the Taylor bubble ceases moving up or down in the axial direction, and hence the pseudo-steady solution is reached, the velocity is then adjusted and set to be the terminal velocity. The initial guess of the liquid film thickness (δ_{LF}) is estimated using [Brown \(1965\)](#) equation, which is given by:

$$\delta_{LF} = \left[\frac{3\nu}{2g(R - \delta_{LF})} U_{TB} (R - \delta_{LF})^2 \right]^{1/3} \quad (10)$$

Referring to [Fig. 1](#), using MRF the inlet flow boundary condition is applied with liquid entering at average uniform velocity equal to velocity of Taylor bubble, $U_{in} = U_{TB}$, $V_{in} = 0$. At the bottom of the domain, the outflow boundary condition is applied as the liquid phase is the only phase available. The symmetry boundary condition is applied at the pipe centreline: $\partial U/\partial y = 0$, $V_{out} = 0$. At the wall, the no-slip condition is applied with wall moving downwards with the following velocities: $U_{wall} = U_{TB}$, $V_{wall} = 0$. The gas phase usually has lower density and viscosity than the liquid phase, thus, full slip can be assumed at the gas-liquid interface. The internal circulatory flow within Taylor bubbles has almost negligible effect on the outer surrounding liquid leading to zero interfacial shear stress at the interface. Thus, the pressure variation within Taylor bubble is small and constant pressure is assumed at the interior of the Taylor bubble ([Akagawa and Sakaguchi, 1966](#); [Morgado et al., 2016](#); [Akagawa and Sakaguchi, 1966](#); [Mao and Dukler, 1990](#);

[Zheng et al., 2007](#); [Morgado et al., 2016](#)). Thus, the kinematic condition, $u \cdot n = 0$, assuming full slip at the gas-liquid interface is applied. The dynamic boundary condition can be divided into two separate boundary conditions: the tangential stress balance assuming zero interfacial shear stress along the interface: $(\bar{\tau} \cdot \hat{n}) \cdot \hat{s} = 0$, and the normal stress balance: $\rho_{il} + \sigma K = \text{constant}$. According to [Mao and Dukler \(1990\)](#), the curvature of the interface (K) is expressed in terms of radii of the curvature of the bubble surface, as follows:

$$K = \left[\frac{1}{r_1} + \frac{1}{r_2} \right] \quad (11)$$

where: r_1 and r_2 are the local principal radii of curvature at the bubble surface as indicated by [Mao and Dukler \(1990\)](#).

3.2. Solution strategy and convergence criterion

A time-dependent simulation is carried out in the present case to investigate the unsteady behaviour of two-phase slug flow. The simulation is carried out using the explicit VOF model. The PISO pressure-velocity is selected. The spatial discretization scheme used is as follows: Green-Gauss Node Based for the gradient, PRESTO for pressure, Geo-reconstruct for volume fraction, a Quick scheme for momentum, and first-order implicit for unsteady formulation. The scaled absolute values of the residual of the calculated values of mass, velocity in x, and y directions are monitored, and convergence criterion of 10^{-4} is set for each time step, with a maximum number of iterations of 1000. The variable time step is applied to the governing equations with global Courant number fixed to 0.25.

4. Results and discussions

In this section, a mesh dependence study is firstly introduced, which is then followed by validation study of single Taylor bubble rising through aqueous glycerol solution and finally the main results. The present study aims at investigating the main hydrodynamic features of the rise of single Taylor bubble through stagnant Newtonian fluid including Taylor bubble shape, Taylor bubble rise velocity (U_{TB}), liquid film thickness (δ_{LF}), maximum liquid film velocity (U_{LF}), wall shear stress (τ_w), and wake length (L_w), and wake volume (V_w) (as shown in [Fig. 1](#)) with particular focus on the governing dimensionless numbers: Eötvös number and Reynolds number. The results are then divided into four sections: Taylor bubble shape, Taylor bubble rise velocity, flow in the liquid film region, and flow in the wake region. In addition, correlations for the terminal velocity of the Taylor bubble, and for the dimensionless wall shear stress are proposed showing the significance of these main dimensionless parameters.

[Table 1](#) lists the simulation cases and their corresponding results. The ranges of Eo , $Re_{U_{TB}}$, $Fr_{U_{TB}}$, Γ_ρ , Γ_μ , $(L_{TB}/D)_i$ are 6–700, 2.6–165, 0.031–0.330, 60–200, 66.7–6674, and 2–10, respectively. The input parameters prior simulation are: Eo , Γ_ρ , Γ_μ , $(L_{TB}/D)_i$, while the predicted values from the simulation are: $Re_{U_{TB}}$, $Fr_{U_{TB}}$, $(L_{TB}/D)_i$, (δ_{LF}/D) , $U_{LF,max}$, $\left(\frac{\tau_w}{\rho_l g D} \right)_{max}$, L_w/D , and V_w/D^3 . The initial ratio of bubble size to pipe diameter is given by $(L_{TB}/D)_p$.

4.1. Mesh dependence study

A uniform grid of quadrilateral control elements is used in the present CFD simulation. A mesh dependence test is done to ensure grid independence results. Three different grid densities are used to simulate a selective experimental case of [Campos and De Carvalho \(1988\)](#) with air and aqueous glycerol solution as working fluids in 19 mm inner diameter vertical pipe. The relevant dimensionless numbers of the selected case are: $Eo = 64$, $Re_{U_{TB}} = 60$, $Fr_{U_{TB}} = 0.3409$.

[Table 2](#) shows the mesh characteristics, the relevant selected

hydrodynamic characteristics to be examined and the corresponding deviations ($Eo = 64$, $Re_{U_{TB}} = 60$, $Fr_{U_{TB}} = 0.3409$). The reference mesh for the deviation calculation is the denser mesh (104×1120 elements). It can be assumed that the results are completely independent particularly when using mesh 2 and mesh 3. Subsequently, based on the results shown in Table 1, the current simulations are conducted using either mesh 2 or mesh 3 depending on pipe geometry.

4.2. Validation of the computational code

In this section, a primary validation of the current numerical code based on the experimental work of Nogueira et al. (2006a,b) on a single Taylor bubble rising through stagnant viscous liquid under laminar flow regime is presented. The experimental condition corresponds to $Eo = 186$, $Re_{U_{TB}} = 37.083$, $Fr_{U_{TB}} = 0.3355$. Fig. 2 shows the numerical results for the flow field including the streamlines and velocity vectors for the selected experimental case using moving reference frame (MRF). From this figure, the flow can be divided into three regions: a, b and c namely: Taylor bubble nose region, falling liquid film region and Taylor bubble wake region (liquid slug zone). In the Taylor bubble nose region, the Taylor bubble moves up with velocity (U_{TB}) due to buoyancy, pushing the liquid sideways where liquid film zone starts to develop. In the falling liquid film region, the liquid moves downwards with velocity (U_{LF}), and decreasing liquid film thickness (δ_{LF}). Once a balance between the gravitational and the friction forces is reached, a constant terminal liquid film velocity and thickness is developed. In the Taylor bubble wake region, the falling liquid film starts to plug into the liquid slug ending with highly mixing zone in the wake structure of the bubble.

Moreover, Fig. 3 gives further validation to the numerical code by showing a comparison between the numerical and experimental of Taylor bubble shape in the nose region, which is an essential feature as discussed by Araújo et al. (2013b). It can be seen that the simulation predicts well the experimental data.

In addition, Fig. 4 shows a direct comparison for the resulting numerical velocity field in three different regions: Taylor bubble nose region, liquid film region, and wake region with the experimental work of Nogueira et al. (2006a,b) in fixed reference frame (FRF). The axial velocity is plotted in dimensionless form (u/U_{TB}) and (x/D) refers to point placed in different axial iso-surfaces according to the region examined. The simulation results showed accepted matching with the experimental results in the different regions examined. Adding to this, the numerical axial velocity profile in the liquid film region shows good matching with both the experimental results and the theoretical velocity profile given by Brown (1965) (equation (12)).

$$u_z = \frac{g}{\nu} \left[\frac{R^2 - r^2}{4} - \frac{(R-r)^2}{2} \ln \frac{R}{r} \right] \quad (12)$$

Furthermore, Table 3 shows the numerical, experimental, and theoretical values for some of the main hydrodynamic characteristics of the experimental validation case and their respective deviations. The value of $Fr_{U_{TB}}$ shows deviation of 7.80% compared to experimental value measured by Nogueira et al. (2006a). It also shows good agreement when compared with experimental and theoretical correlations of Wallis (1969) (equation (9)), and Viana et al. (2003) whose theoretical correlation is given by:

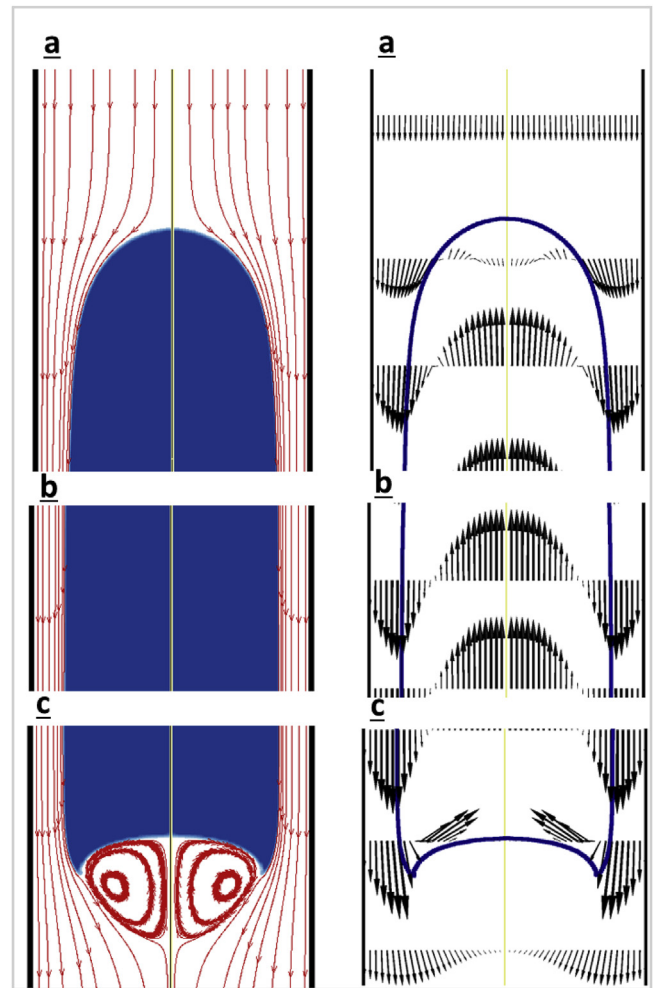


Fig. 2. Numerical results of the Streamlines (left) and Velocity fields (right) of a rising Taylor bubble through stagnant liquid vertical pipe, $Eo = 186$, $Re_{U_{TB}} = 37.083$, $Fr_{U_{TB}} = 0.3355$ using moving reference frame (MRF).

$$Fr = \frac{0.34 / \left(1 + \frac{3805}{Eo^{30.6}} \right)^{0.58}}{\left(1 + \left(\left(\frac{R_G}{31.08} \right) \left(1 + \frac{778.76}{Eo^{1.96}} \right)^{-0.49} \right)^{-1.45} \left(1 + \frac{7.22 \times 10^{13}}{Eo^{9.93}} \right)^{0.094} \right)^{0.74} \left(1 + \frac{7.22 \times 10^{13}}{Eo^{9.93}} \right)^{-0.094}} \quad (13)$$

where $R_G = \sqrt{D^3 g (\rho_L - \rho_G) \rho_L / \mu_L}$.

For the flow in liquid film region, three important parameters are discussed: liquid film thickness (δ_{LF}), wall shear stress (τ_w), and maximum liquid film velocity (U_{LF}). Firstly, the values of δ_{LF} , and τ_w shows good matching with both the experimental data of Nogueira et al. (2006a) and the theoretical predictions of Brown (1965). Brown (1965) equation for δ_{LF} prediction is given by equation (10), while the prediction

Table 2
Mesh dependence test results, $Eo = 64$, $Re_{U_{TB}} = 60$, $Fr_{U_{TB}} = 0.3409$.

Mesh	U_{TB} (m/s)	Error U_{TB} (%)	δ_{LF} (m)	Error δ_{LF} (%)	τ_w (Pa)	Error τ_w (%)	L_w/D	Error L_w/D (%)	V_w/D^3	Error V_w/D^3 (%)
26*280	0.1431	5.03	0.001972	1.72	20.74	1.06	0.5080	8.35	0.781	-5.38
52*560	0.1373	1.02	0.001947	0.46	20.58	0.29	0.480	3.04	0.805	-2.24
104*1120	0.1359	-	0.001938	-	20.52	-	0.4654	-	0.823	-

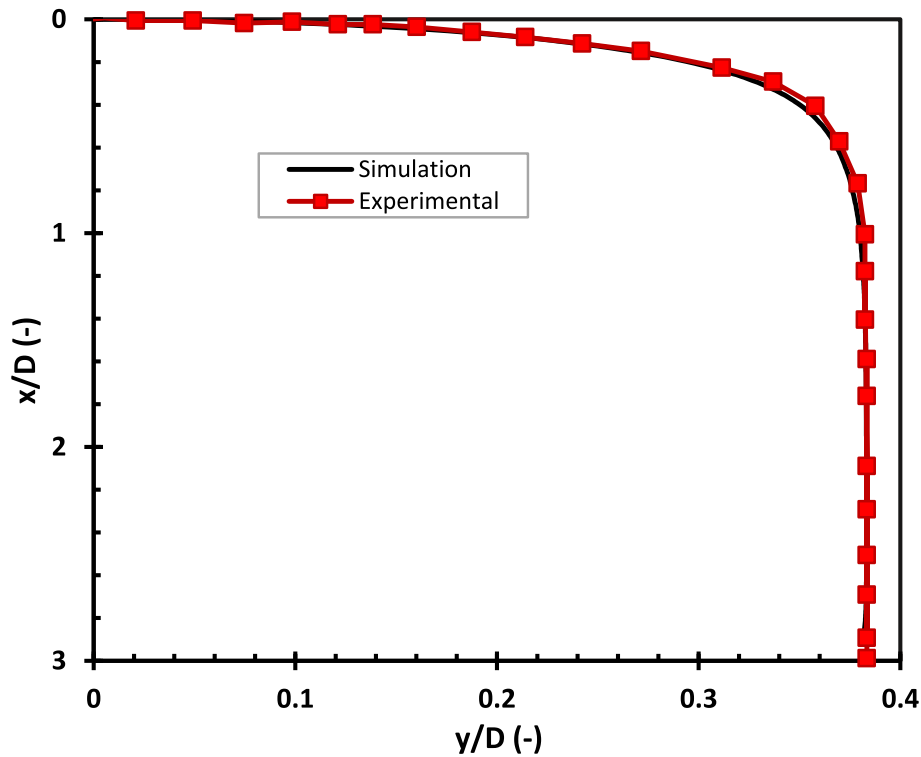


Fig. 3. The experimental and numerical shape of Taylor bubble in nose region - x is the axial distance from bubble nose.

for τ_w is given by:

$$\tau_w = \rho_L g \left[\frac{R}{2} - \frac{(R - \delta_{LF})^2}{2R} \right] \quad (14)$$

Secondly, the value of the third parameter (U_{LF}) is as well close to experimental value. Furthermore, for the flow in the wake region, two main parameters namely: L_w/D , and V_w/D^3 are used to characterize the wake structure. The deviation with experimental values is quite large between 42% and 48% which is clarified by Araújo et al. (2013b) by the fact that the experimental technique did not account for the wake in the concave bottom at the rear of Taylor bubble. However, smaller deviations are established when comparing with the experimental correlations developed by Campos and De Carvalho (1988), which is given by:

$$\begin{aligned} \frac{L_w}{D} &= 0.30 + 1.22 \times 10^{-3} N_f \text{ for } 100 < N_f < 500 \\ \frac{V_w}{D^3} &= 7.5 \times 10^{-4} N_f \text{ for } 100 < N_f < 500 \end{aligned} \quad (15)$$

The present study investigates the effects of density ratio (Γ_ρ), viscosity ratio (Γ_μ), and the bubble size to the pipe diameter (L_{TB}/D) on the hydrodynamic characteristics of the rise of single Taylor bubble in stagnant liquid to support previous numerical work and experimental work in literature. Cases 1 to 3 in Table 1 clearly emphasises that the density ratio has almost no effect on the dynamics of Taylor bubbles. The density ratio has minimal effect on liquid film thickness, maximum the velocity of liquid film, maximum wall shear stress, wake length and wake volume. The simulation results as well agree with the numerical result of Kang et al. (2010). For the viscosity ratio effect, three cases denoted by cases 3, 4, and 5 are simulated with viscosity ratios of 66.7, 667, and 6674, respectively. Referring to Table 1, it is also concluded that the viscosity ratio (Γ_μ) has minimal effect on the dynamics of Taylor bubbles including as well the liquid film thickness, maximum velocity of liquid film, maximum wall shear stress, wake length and wake volume. In conclusion, the simulated cases 1 to 6 have almost the same values of $Re_{U_{TB}}$, and Eo , which is 24.34, and 66, respectively. Thus, the simulation

results further assist the conclusion made by Lu and Prosperetti (2008) and Kang et al. (2010) that the both density ratio (Γ_ρ), and viscosity ratio (Γ_μ) have a negligible effect on the dynamics of Taylor bubbles. Lastly, to explore the effect of L_{TB}/D , four cases are simulated denoted by cases 7 to 10 in Table 1 corresponding to L_{TB}/D of 2, 3, 4, and 10, respectively. It can be also concluded that L_{TB}/D as well has minimal effect on the dynamics of Taylor bubbles. Table 1 shows that these four cases has almost the same value of $Fr_{U_{TB}}$ which agrees well with the conclusion of Kawaji et al. (1997) that the bubble length has no effect on the bubble terminal speed and subsequently no effect on Froude number ($Fr_{U_{TB}}$).

Further conclusion indicates that the validation results, in addition to the validation given by Massoud et al. (2016), where four cases based on the experimental work of Campos and De Carvalho (1988) are simulated and show good agreement when compared with both experimental and numerical data, illustrate that the numerical code is capable of simulation of Taylor bubble rising through stagnant liquid with high satisfaction degree of results. The problem can now be treated in terms of three main dimensionless groups, given by:

$$Fr_{U_{TB}} = \frac{U_{TB}}{\sqrt{gD}} = f \left[Eo = \frac{g\rho_L D^2}{\sigma}, Re_{U_{TB}} = \frac{\rho_L U_{TB} D}{\mu_L} \right] \quad (16)$$

In the following section, the effect of these core dimensionless groups that govern the present problem with the guidelines of the order of magnitude analysis are examined on the developed Taylor bubble shape, Taylor bubble rises velocity, flow in the liquid film and flow into the wake region.

4.3. Taylor bubble shape

In this section, the effect of the main dimensionless groups that govern the present problem given by relation (16) is discussed. This includes representing the final shape of the developed Taylor bubble, the flow field around it (streamlines), and the Taylor bubble shape profile for each of the dimensionless groups examined.

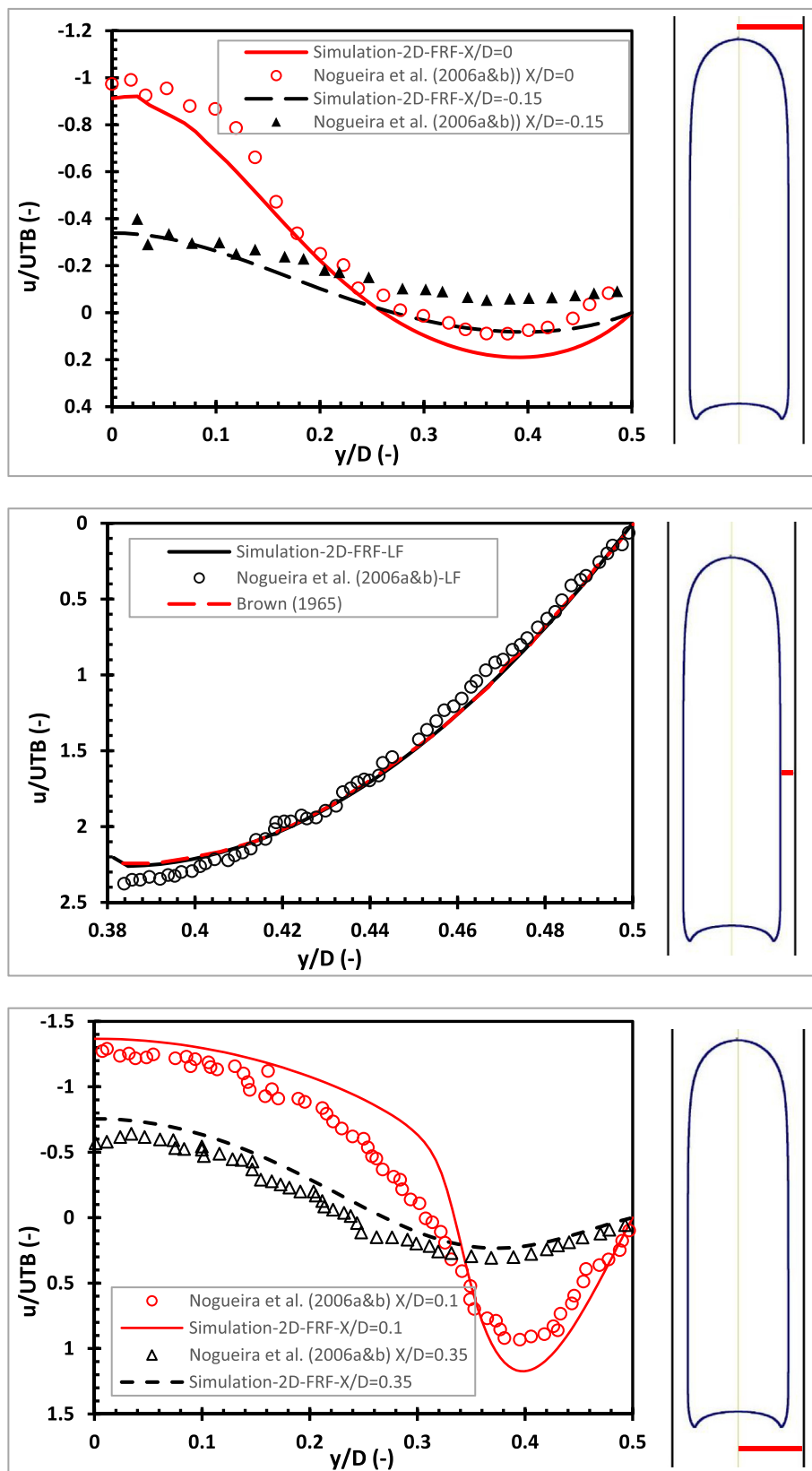


Fig. 4. Numerical, experimental, and theoretical dimensionless axial velocity profile for three different regions: Taylor bubble nose region, liquid film region, and wake region in a fixed reference frame. $Eo = 186$, $Re_{UTB} = 37.083$, $Fr_{UTB} = 0.3355$.

Table 3

Numerical, experimental, and theoretical values for some of the main hydrodynamic characteristics of experimental case based on work of Nogueira et al. (2006a,b) with $Eo = 186$, $Re_{U_{TB}} = 37.083$, $Fr_{U_{TB}} = 0.3355$, and their respective deviations.

Simulation		Experimental		Theoretical
1. Taylor bubble rise velocity-Froude number				
$Fr_{U_{TB}}$ (-)	0.3094	0.3355 ^a	0.3055 ^b	0.3309 ^c
Error (%)		7.80	-1.26	6.50
2. Flow in the liquid film region				
δ_{LF} (m)	0.0037	0.00384 ^a	0.00382 ^a	
Error (%)		2.73	2.37	
τ_w (Pa)	39.83	39.6 ^a	40.75 ^d	
Error (%)		-0.59	2.22	
U_{LF} (m/s)	0.279	0.253 ^a	-	
Error (%)		-10.02	-	
3. Flow in the wake region				
L_w/D (-)	0.371	0.19 ^a	0.4347 ^e	-
Error (%)		48.78	-17.19	-
V_w/D^3 (-)	0.069	0.04 ^a	0.0828 ^e	-
Error (%)		42.67	-18.49	-

^a Nogueira et al. (2006a,b).

^b Viana et al. (2003).

^c Wallis (1969).

^d Brown (1965).

^e Campos and De Carvalho (1988).

4.3.1. Effect of Reynolds number

In literature, most of the studies done on dynamics of Taylor bubble rising through stagnant liquid highlight the major effect of inverse viscosity number or Archimedes number without considering the significant effect of Reynolds number. In this section ten cases namely: cases 11 to 20 in Table 1 are simulated to investigate the significant effect of $Re_{U_{TB}}$ on the dynamics of Taylor bubble. Fig. 5 demonstrates the effect of $Re_{U_{TB}}$ number on the final shape of the developed Taylor bubble and the flow field for cases 11, 13, 15, and 18 with $Re_{U_{TB}}$ values corresponding to 12, 46, 80, and 131, respectively. It can be shown that $Re_{U_{TB}}$ has prodigious

effect on the final shape of bubble, as indicated in Fig. 5. Particularly, for low values of $Re_{U_{TB}}$, the viscous forces of liquid phase surrounding the bubbles are high enough to encumber the rise of the Taylor bubble, and hence the terminal bubble velocity will be at its lowest values (lowest $Fr_{U_{TB}}$), as indicated in Table 1. The increase in $Re_{U_{TB}}$ increases the velocity of Taylor bubble. It can be concluded from Table 1 that for $Re_{U_{TB}}$ values approximately above 80, $Fr_{U_{TB}}$ is almost the same which is in agreement with the experimental conclusions of Wallis (1969). It can also be concluded from Fig. 5, that the high viscous forces enhance the elongation of the terminal developed Taylor bubble. The gradual increase in $Re_{U_{TB}}$ increases the concavity of the rear of Taylor bubble which is in good agreement with the experimental observation of Goldsmith and Mason (1962) that the rear of Taylor bubbles is characterized by being flat or concave when the flow is not viscosity dominated, and oblate spheroid when it is viscosity dominated. The shape of the rear of Taylor bubble and its transition critical criteria is discussed in details in section 4.6.

It is also concluded from Fig. 5, that $Re_{U_{TB}}$ has significant effect on wake development. For low values of $Re_{U_{TB}}$, the liquid from the liquid film region expands directly and smoothly over the Taylor bubble bottom which is noticed by parallel streamlines in the wake region. Increasing the values of $Re_{U_{TB}}$, the Taylor bubble becomes wider in size squeezing the liquid film region that subsequently increases the velocity of the trailing liquid plugging in to the Taylor bubble bottom. This leads to the development of circulatory rear of vortices as liquid plugs into the Taylor bubble bottom. The scale and intensity of the vortex increases with higher values of $Re_{U_{TB}}$. In addition, Fig. 6 shows the effect of $Re_{U_{TB}}$ on the bubble shape profile where it is clearly seen that the increase in $Re_{U_{TB}}$ causes the bubble nose to become less slender, and reduces the thickness of the developed liquid film.

4.3.2. Effect of Eötvös number

Eötvös number (Eo) represents the effect of buoyancy and surface tension forces which are two of the significant forces acting on Taylor bubbles. In order to investigate the effect of Eo number on the dynamics of Taylor bubble rising through stagnant liquids, 21 cases namely case 21 to case 42 in Table 1 are simulated. This covers wide range of Eo varying from 6 to 700.

Fig. 7 shows the effect of Eo on the final shape of the developed Taylor bubble and the flow field for cases 22, 25, 28 and 32 with Eo values corresponding to 10, 66, 100 and 160, respectively. At low values of Eo , the surface tension forces are high enough to maintain any distortion at the gas-liquid interface. The bubble is noticed to have prolate spheroidal

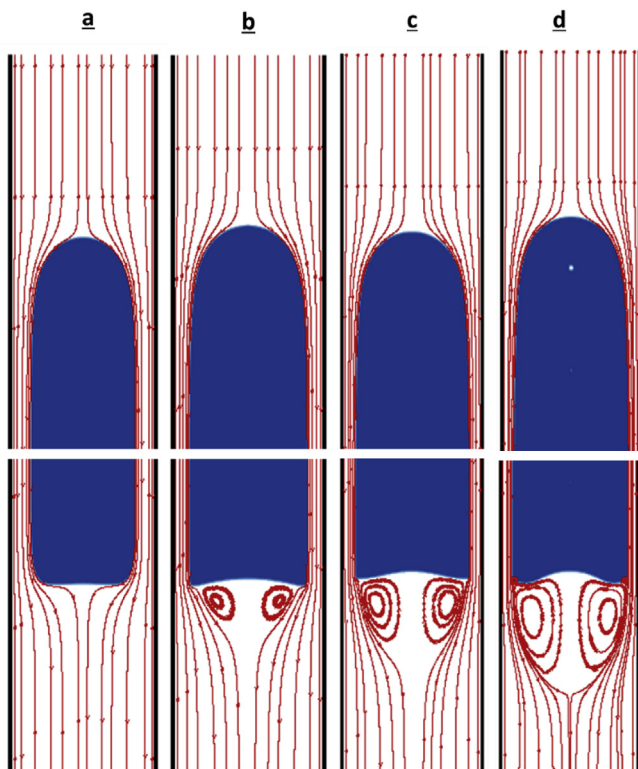


Fig. 5. Effect of $Re_{U_{TB}}$ on the terminal shape of Taylor bubble, and streamlines representing the flow field. (a) Case (11) $Re_{U_{TB}} = 12$, (b) Case (13) $Re_{U_{TB}} = 46$, (c) Case (15) $Re_{U_{TB}} = 80$, and (d) Case (18) $Re_{U_{TB}} = 131$.

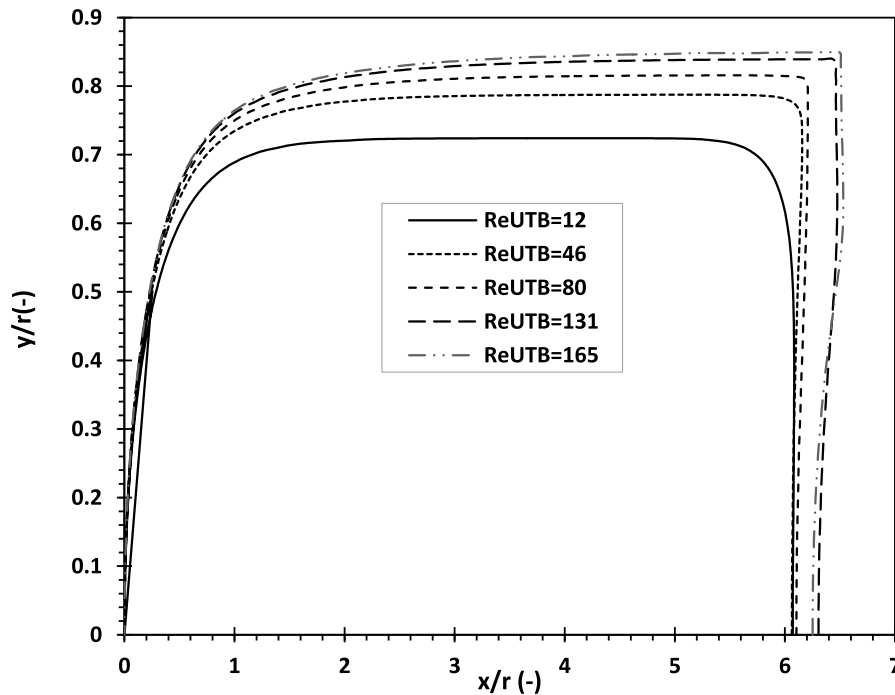


Fig. 6. Effect of Re_{UTB} on bubble shape profile - x is the axial distance from bubble nose.

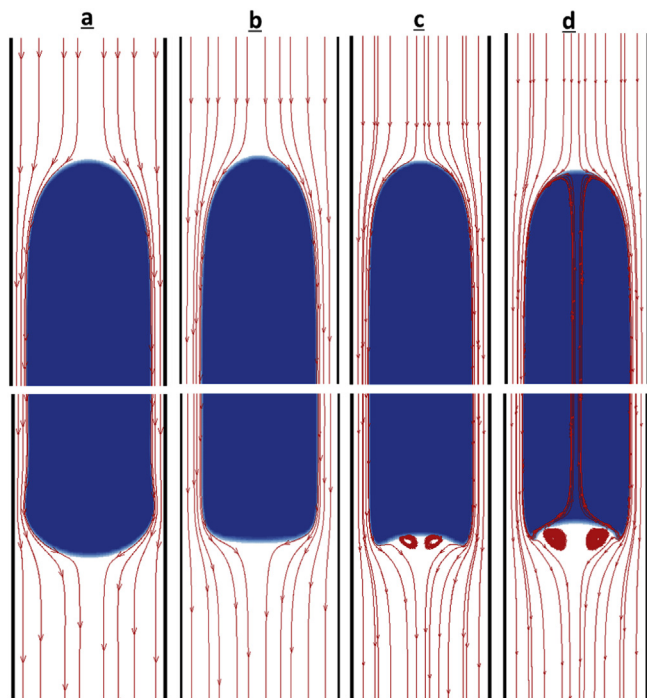


Fig. 7. Effect of Eo on the terminal shape of Taylor bubble, and streamlines representing the flow field. (a) Case (22) $Eo = 10$, (b) Case (25) $Eo = 66$, (c) Case (28) $Eo = 100$, and (d) Case (32) $Eo = 160$.

nose and oblate spheroid bottom. The increment in the surface tension increases the curvature of the bubble nose, and subsequently increases the liquid film thickness. Hence, the velocity of fully developed falling liquid film decreases.

Regarding the rear of the bubble, the increase in Eo leads to increase in inertial forces which is seen as values of Fr_{UTB} increases. The increase in Eo also significantly affects the rear of bubble by gradually turning the

bottom shape from convex or flat into concave structure, as indicated in Fig. 7. This causes wakes to be developed at the rear of the bubble, and hence wake length and volume increases with the increase in Eo values, as indicated in Table 1. This has similar trend as that shown in Fig. 5 exploring the effect of Re_{UTB} . At certain critical value of Eo , the further increase in Eo with its corresponding deficiency of surface tension results in deformation at gas-liquid interface. The shape of the bubble's rear turns from concave into wave and subsequently into breaking up. The breaking up concept is characterized by very unstable phenomena that occurs at the rear of the bubble and is noticed by the small bubbles shedding off from Taylor bubble into the wake region. Higher values of Eo , cause the gas-liquid interface to lose its structure and strong breaking up concept is noticed (Morgado et al., 2016). A phase diagram map is discussed later in section 4.6 that describes the interesting phenomena that occur at the rear of Taylor bubbles which is strongly affected by surface tension, and inertia forces.

Furthermore, Fig. 8 shows the effect of Eo on the bubble shape profile for different selected cases from Table 1. The increase in Eo increases the bluntness of the nose of the bubble, decreases the flatness of the bubble tail, and increment the liquid film thickness. This conclusion refutes that of Kang et al. (2010) that Eo has no effect on the dynamics at bubble nose, and liquid film thickness, but agrees with the conclusion of Taha and Cui (2006) and Zheng et al. (2007) which entails that Eo increases the degree of prolateness of the Taylor bubble nose.

In conclusion, Eo , and Re_{UTB} significantly control the developed shape of Taylor bubble. The results shows that the surface tension forces are significant with low values of Eo , while the large contribution of Re_{UTB} on dynamics of Taylor bubble is noticed with higher values of Re_{UTB} . These conclusions agree well with the results of order of magnitude analysis discussed earlier in section 2. In the following section, the role of these dimensionless groups on the Taylor bubble rise velocity is introduced.

4.4. Taylor bubble rise velocity

The Taylor bubble rise velocity (U_{TB}) is one of the main hydrodynamic features used for the description of two phase slug flow. In this section, a detailed discussion on Taylor bubble terminal velocity is

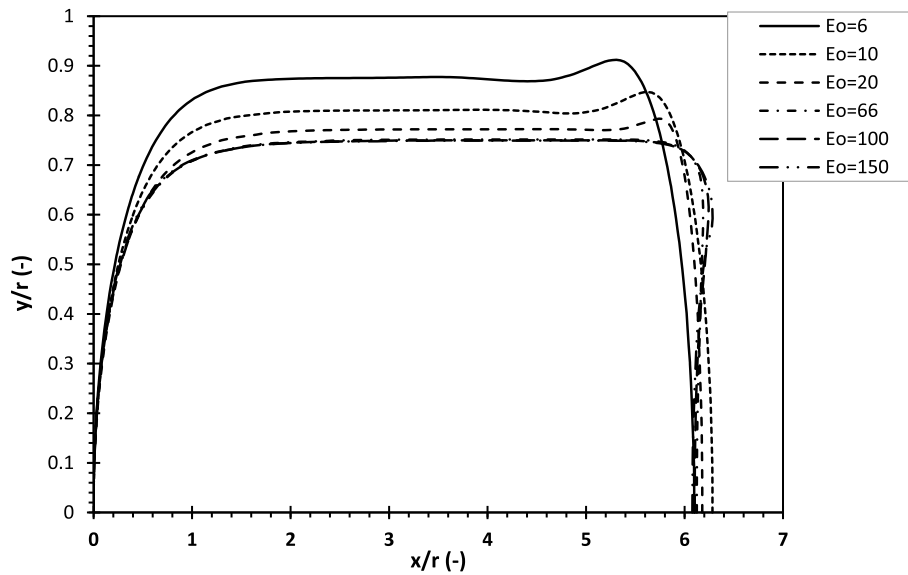


Fig. 8. Effect of Eo on bubble shape profile - x is the axial distance from bubble nose.

introduced. Based on the discussion given in section 2, it can be concluded that the terminal bubble velocity (U_{TB}) is mainly governed by viscosity, surface tension, buoyancy and inertia forces. In literature, various studies either theoretical and/or experimental are done to account for the terminal bubble velocity. A good review for the main correlations for $Fr_{U_{TB}}$, to account for bubble terminal velocity, starting from the theoretical investigation of Dumitrescu (1943) and ending with the recent study of Kurimoto et al. (2013) and is given by Morgado et al. (2016).

Based on relation (16), and on the cases simulated in Table 1, a proposed correlation to estimate $Fr_{U_{TB}}$ that depends mainly on Eo and $Re_{U_{TB}}$ is developed, given by:

$$Fr = \frac{U_{TB}}{\sqrt{gD}} = 0.0359 - 0.3596Eo - 0.7067Re_{U_{TB}} + 0.5801Eo^2 - 1.014EoRe_{U_{TB}} + 0.3447Eo^3 + 1.594Eo^2Re_{U_{TB}} - 0.1931Eo^4 + 0.9647Eo^3Re_{U_{TB}} - 0.001814Eo^5 - 0.5481Eo^4Re_{U_{TB}} \quad (17)$$

Additionally, Fig. 9 shows that the values obtained from the proposed correlation fit quite well with the behaviour of other well-known correlations from literature for different domains. Values predicted from Viana et al. (2003) correlations, and values from the recent correlation of Kurimoto et al. (2013) are added to Fig. 9 for comparison purposes. It can be seen that the proposed correlation agrees well with a wide range of correlations with a maximum deviation of $\pm 10\%$. In general, there is a proper matching between the data, especially in surface tension dominant domain. However, in inertia dominant domains with high values of

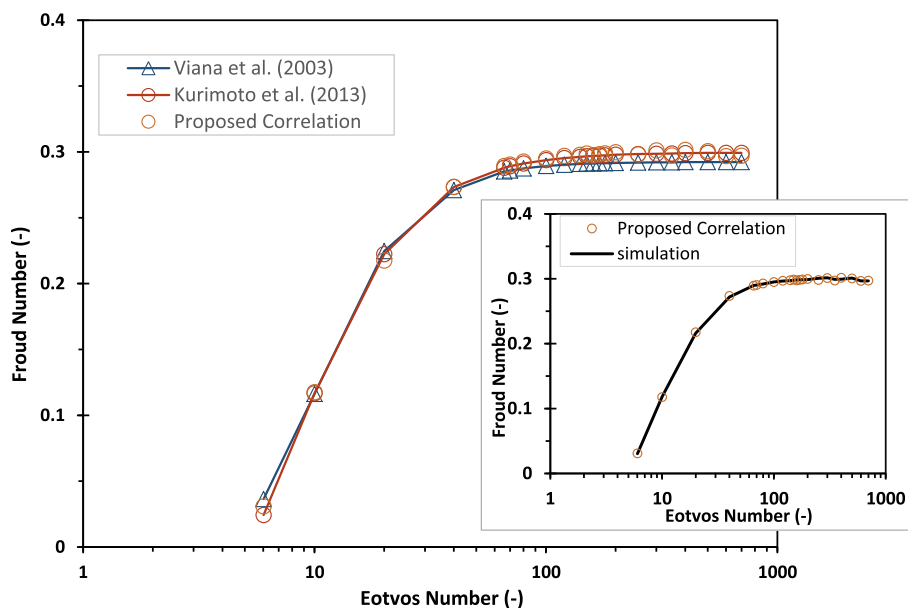


Fig. 9. Numerical results of U_{TB} expressed in terms of $Fr_{U_{TB}}$ for a several Eo with corresponding values obtained from proposed correlation given in equation (17), and different correlations from literature for comparison. $Re_{U_{TB}}$ varies from 2.6 to 25.3.

Eo , the proposed correlation predicts the correlation of Viana et al. (2003). The only explanation is most probable to be a numerical problem for situations of high inertial forces as pointed out by Araújo et al. (2012). In brief, the proposed correlation shows an accepted behaviour with other correlations with maximum deviation $\pm 10\%$.

4.5. Flow in liquid film

The flow in the liquid film is investigated by characteristics associated with three key features: the dimensionless developed liquid film thickness (δ_{LF}/D), the dimensionless velocity of the liquid film (U_{LF}/U_{TB}) and the dimensionless wall shear stress $\left(\frac{\tau_w}{\rho_l g D}\right)_{max}$. This section will discuss the effect of the main dimensionless parameters, given in section 2, on these three key features characterizing the flow in the liquid film.

4.5.1. Liquid film velocity and thickness

Fig. 10 represents the effect of Re_{UTB} on the dimensionless normalized thicknesses and velocity of falling liquid film along the Taylor bubble length. At low values of Re_{UTB} , the liquid film thickness (δ_{LF}/D) decreases with the increase in the dimensionless distance measure from bubble nose (x/D) until it reaches a constant thickness at around $x/D = 1$. At that point, a balance between the gravitational and friction forces is reached, and hence a constant liquid film thickness and velocity is established. However, the dimensionless velocity of the liquid film (U_{LF}/U_{TB}) changes contrarily. The increase in Re_{UTB} diminishes the long slender shape of Taylor bubble and turns it into shorter and flatter bubbles which in turns reduces (δ_{LF}/D) (Kang et al., 2010). In addition, the higher viscosity of the surrounding liquid at low values of Re_{UTB} increment the shear stress in the liquid, thus, the momentum diffusion opposes the liquid flowing from the liquid film into the bubble wake region (Zheng et al., 2007). Therefore, the increase in Re_{UTB} increases the dimensionless velocity of the liquid film (U_{LF}/U_{TB}).

Finally, the effect of Eo on the dimensionless normalized thicknesses and velocity of falling liquid film along the Taylor bubble length is given in Fig. 11. As discussed earlier, the increase Eo affects the curvature of the bubble nose, thus increases (δ_{LF}/D), increases stabilization length (distance needed for formation of fully developed falling liquid film), and finally increases U_{LF} . However, the dimensionless velocity of the liquid

film (U_{LF}/U_{TB}) decreases with the increment in Eo due to the fact that the inertia forces are increased leading to increment in U_{TB} values. It is concluded from Fig. 11 that (δ_{LF}/D) and (U_{LF}/U_{TB}) are affected with the increase in Eo till $Eo \leq 66$, where further increase in Eo shows almost no effect on the flow in the liquid film.

In conclusion, it should be pointed out that both of (δ_{LF}/D), and (U_{LF}/U_{TB}) are strongly dependent on Eo , and Re_{UTB} . This conclusion contradicts those made by: Kang et al. (2010): the liquid film thickness is only dependent on Archimedes number ($Ar = N_f^2$), and that of Araújo et al. (2012): the liquid film thickness is function of N_f , and M .

4.5.2. Wall shear stress distribution

If the two-phase slug flow problem is involved in heat or mass transfer, then the wall shear stress becomes a primary significant hydrodynamic parameter. This process is often referred to as slug flow-induced corrosion (Zheng and Che, 2006; Zheng and Che, 2007; Zheng et al., 2007; Araújo et al., 2012). The main problems result from slug flow corrosion are: pipeline damage, decrease pipe lifetime and may lead to the shutdown of the pipeline.

Fig. 12 shows the wall shear stress distribution along the Taylor bubble length for different Re_{UTB} . For low values of Re_{UTB} , the wall shear stress distribution starts with an increase in the wall shear stress near the bubble nose then it reaches a maximum positive value with the formation of a constant liquid film characteristics (thickness and velocity). The wall shear stress then starts to decrease till it reaches zero value in the bubble tail or wake region.

The increase in Re_{UTB} , results in less the viscous liquid surrounding the Taylor bubble, that subscribes to decrement in wall shear stress. This conclusion agrees well with that made by Taha and Cui (2006). On the one hand, the dimensionless wall shear stress in the nose region is not affected by the increase in Re_{UTB} , on the other hand, the plateau behaviour at the developed liquid film is shortened with the increase in Re_{UTB} . This occurs as a result of the shape in Taylor bubble shape which is characterized by being a long slender that turns into shorter and flatter bubbles with the increase in Re_{UTB} . Additionally, it should be pointed out that the effect of increase in Re_{UTB} on the bubble wake region is seen as jump in dimensionless wall shear values that increases with higher values of Re_{UTB} ; this further assists the conclusion made by Kang et al. (2010). A further validation for the flow in liquid film region is given in Fig. 13,

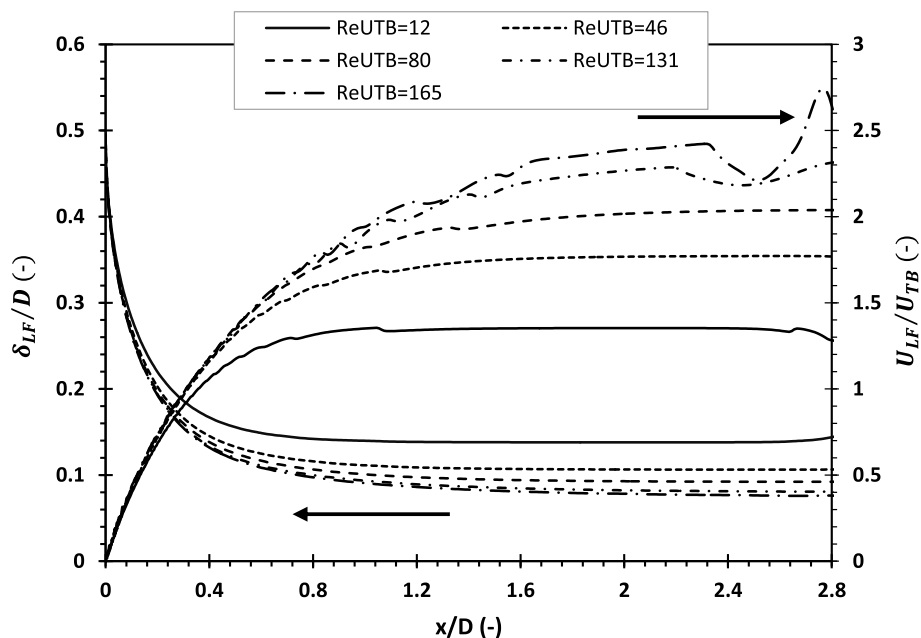


Fig. 10. The effect of Re_{UTB} on the dimensionless normalized thicknesses (δ_{LF}/D) and dimensionless velocity of falling liquid film (U_{LF}/U_{TB}) along the Taylor bubble length - x is axial distance from bubble nose.

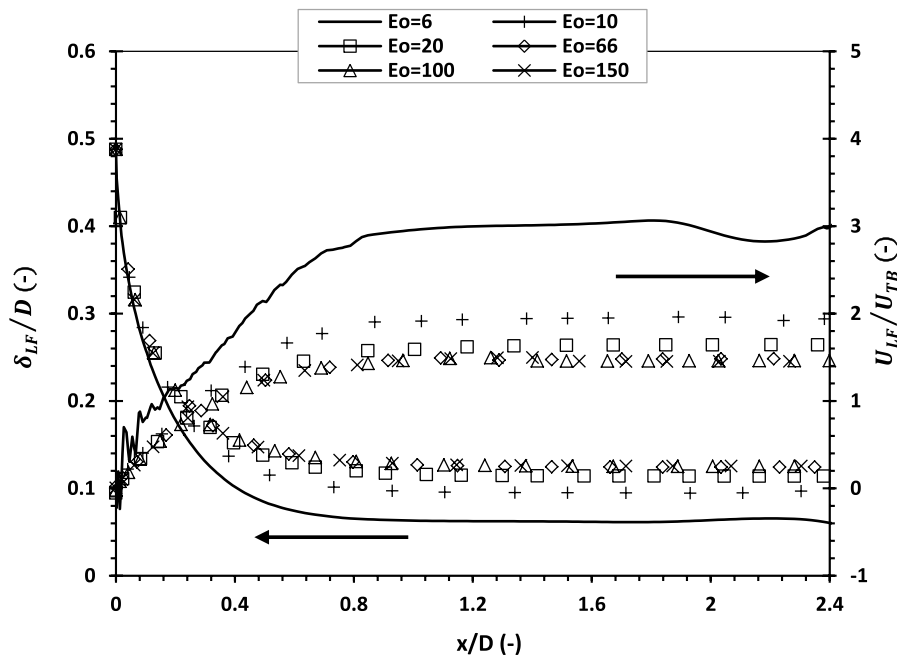


Fig. 11. The effect of Eo on the dimensionless normalized thicknesses (δ_{LF}/D) and dimensionless velocity of falling liquid film (U_{LF}/U_{TB}) along the Taylor bubble length - x is axial distance from bubble nose.

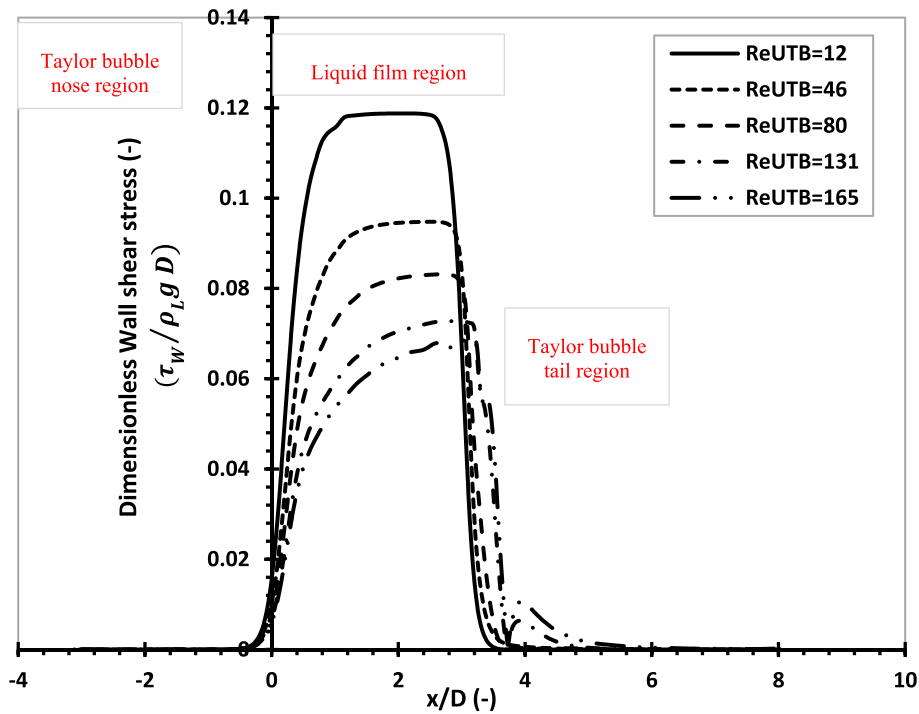


Fig. 12. Effect of Re_{UTB} on the wall shear stress distribution $\left(\tau_w / \rho_L g D\right)_{max}$ along Taylor bubble length - x is axial distance from bubble nose.

where a comparison takes place between the numerical results for dimensionless wall shear stress, and δ_{LF}/D , and the theoretical values by Brown (1965). It is clear that the numerical results strongly correspond with the theoretical data.

Finally, the effect of Eo on the distribution of the dimensionless wall shear stress along the Taylor bubble is given in Fig. 14. Generally, for most values of Eo , the distribution is almost the same, which starts with increase in wall shear stress near the bubble nose, then a constant value is reached at the developed liquid film, followed by reduction in dimensionless wall shear stress near the bubble tail. As indicated in Fig. 14,

lower values of Eo possess different distribution with noticeable peak value in dimensionless shear stress right before the Taylor bubble rear. This is due to the sharp flat and convex shape of the bubble attained at lower values of Eo as discussed in section 4.3.2 which coincides with Araújo et al. (2012).

In conclusion, the numerical results show that the dimensionless wall shear stress is dependent on both Eo , and Re_{UTB} which again contradicts the conclusion made by Kang et al. (2010) that the wall shear stress is only function of Ar , and supports Araújo et al. (2012) conclusion in different scenario. Subsequently, a correlation based on the numerical

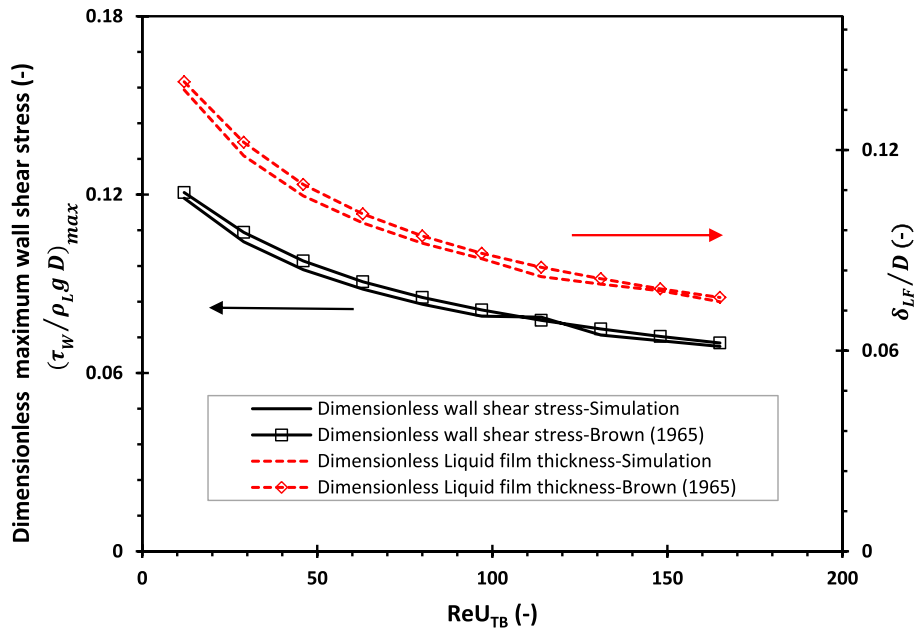


Fig. 13. Variation of the dimensionless normalized thickness (δ_{LF}/D) and dimensionless maximum wall shear stress $\left(\tau_w / \rho_L g D\right)_{max}$ distribution along Taylor bubble length with $Re_{U_{TB}}$ plotted with the theoretical prediction of Brown (1965) (equation (10), and equation (14)).

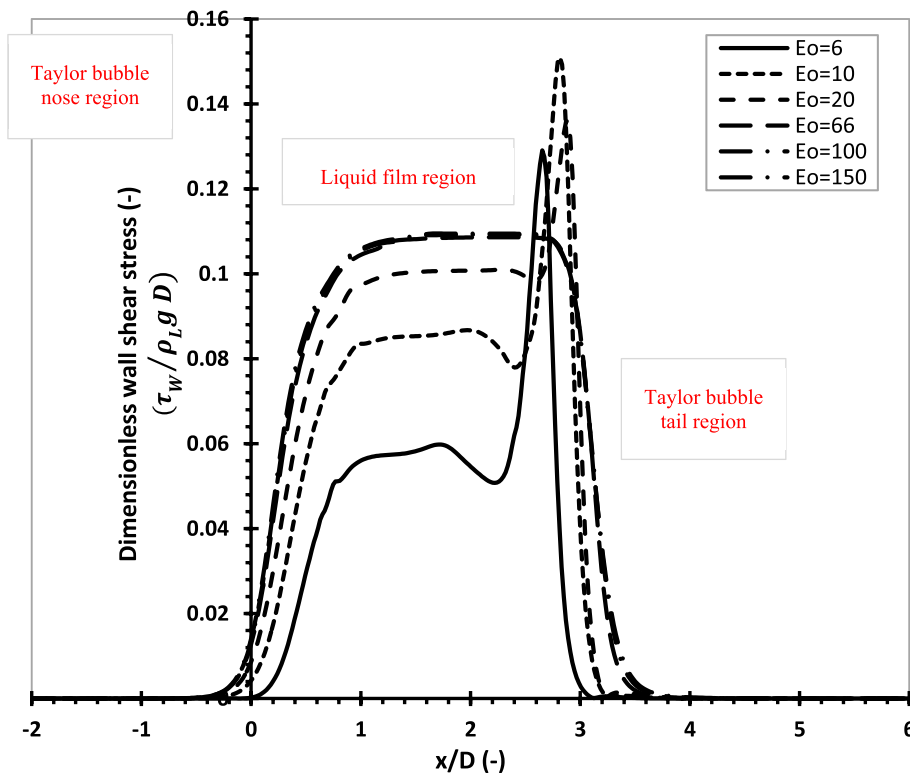


Fig. 14. Effect of Eo on the wall shear stress distribution along Taylor bubble length - x is axial distance from bubble nose, $Re_{U_{TB}}$ varies from 2.6 to 25.

results for all simulated cases in Table 1 for dimensionless maximum wall shear stress is suggested, given by:

$$\begin{aligned} \left(\frac{\tau_w}{\rho_L g D}\right)_{max} = & 0.242 + 0.238Eo + 0.5544Re_{U_{TB}} - 0.1196Eo^2 + 1.037EoRe_{U_{TB}} \\ & + 0.5294Re_{U_{TB}}^2 - 2.818e^{-05}Eo^3 - 0.3202Eo^2Re_{U_{TB}} \\ & + 1.083EoRe_{U_{TB}}^2 - 0.0008579Re_{U_{TB}}^3 \end{aligned} \quad (18)$$

Fig. 15 illustrates the present numerical results, and values obtained from proposed correlation given in equation (18) for the maximum dimensionless wall shear stress along the Taylor bubble length for different values of Eo (cases 21 to 42 in Table 1). It is obvious that the correlation matches the simulation data to a considerable extent. For comparison issues, the theoretical prediction given by Brown (1965) (equation (14)), and predictions using Kang et al. (2010) correlation (equation (19)) are added to Fig. 15.

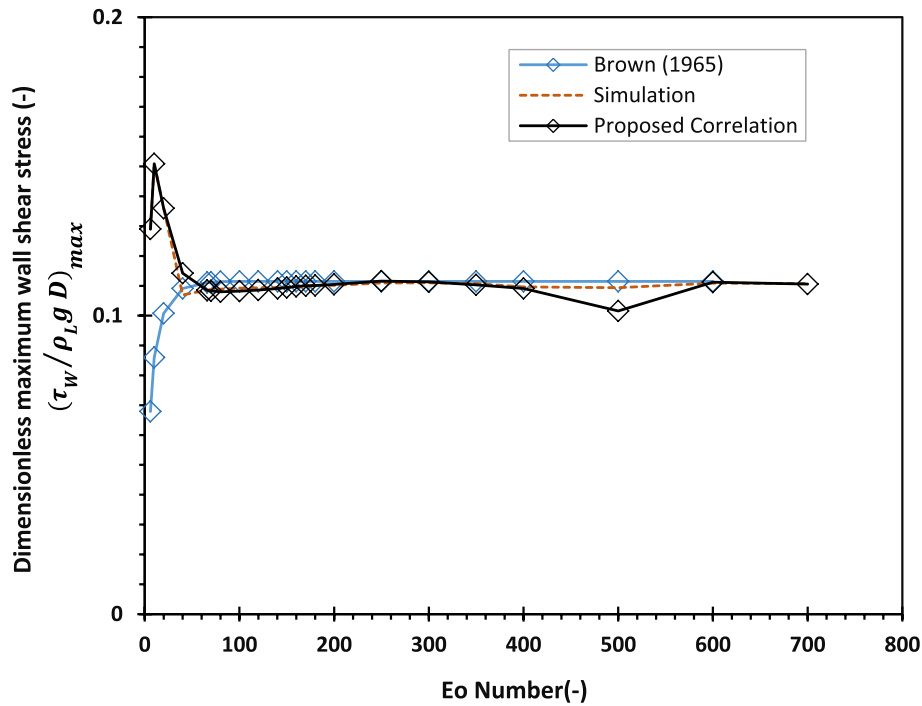


Fig. 15. Numerical results of $\left(\tau_w / \rho_L g D\right)_{max}$ for a several Eo with corresponding values obtained from proposed correlation given in equation (18), theoretical prediction given by Brown (1965) (equation (14)), and predictions using Kang et al. (2010) correlation (equation (19)), for cases 21 to 42 in Table 1.

$$\left(\frac{\tau_w}{\rho_L g D}\right)_{max} = -0.02 \log_{10} Ar + 0.2 \quad (19)$$

It can be seen from Fig. 15, that the suggested correlation matches well with the theoretical predictions calculated by Brown (1965) with very small deviations. However, there is an exception for that proper matching for small values of Eo . Generally, these cases with low values of Eo (especially around 6) possess low values of U_{TB} with almost non-moving bubbles. These cases are more sensitive to numerical errors as clarified by (Zheng et al., 2007; Araújo et al., 2012).

It should be pointed out that $Re_{U_{TB}}$ is not constant for these selected cases, and its values are indicated in Table 1. Similarly, as discussed before for the liquid film thickness, the dimensionless wall shear stress is function of both Eo , and $Re_{U_{TB}}$, and not only function of Ar as concluded by Kang et al. (2010). This conclusion completely agrees with Araújo et al. (2012) that dimensionless wall shear stress depends on Eo , and M , but with different scenario as discussed in section 2.

In conclusion, the flow in the liquid film is considerably affected by Eo , and $Re_{U_{TB}}$. It should be pointed out that both Eo , and $Re_{U_{TB}}$ significantly control the wall shear stress distribution which in turns control the corrosion process related to two-phase flow which is known as slug flow induced corrosion. In the next section, the contribution of these dimensionless groups on the flow in the wake region is discussed.

4.6. Flow in wake region

The wake structure is one of the vital hydrodynamic characteristics of slug flow systems, especially in describing the interaction and coalescences between successive Taylor bubbles and in modelling process (Araújo et al., 2012). Understanding the mechanism by which the wake region is developed is essential prior to introducing the simulation results for flow in wake region. The annular falling film from the liquid film region plugs into the rear of the Taylor bubble and creates mixing and recirculation zone which is known by bubble wake. The intensity and size of the recirculation vortices in the wake region depend on the properties of the surrounding liquid, especially its viscosity as indicated by the

experimental work of Campos and De Carvalho (1988). As discussed previously in section 1, Campos and De Carvalho (1988) conclude three flow patterns for the wake depending on the inverse viscosity number, N_f . In the present study, most of the cases bump into closed axisymmetric laminar wake.

Stimulated by the work of Morgado et al. (2016), and based on the simulated cases given in Table 1, phase diagram of the presence and nature of wake, and of the shape of the rear Taylor bubble are presented in Fig. 16, and Fig. 17, respectively. The diagrams are given in terms of Eo , and $Re_{U_{TB}}$, based on the problem formulation given in section 2. Throughout the present study, only two kinds of wake are observed: either closed axisymmetric wake or no wake. For the shape of rear of Taylor bubble, as discussed earlier in section 4.3.2, when the surface tension is significantly reduced, the gas-liquid interface is easily deformed and the shape of Taylor bubble rear becomes unstable. Stable bubble shape is classified in to concave, convex, or flat bubbles. However, unstable bubble wakes are either wavy or breaking up bubbles. Fig. 17 showed the type of bubble rear shape for all data simulated which strongly corresponds with Kang et al. (2010) and Morgado et al. (2016) which conclude that the sudden elongation in the bubble tail, based on investigating the effect of Eo on the dynamics of Taylor bubbles, is around unity which corresponds to $Eo > 250$. Referring to cases 21 to 42 in Table 1, it is observed that the gas-liquid interface starts to become unstable approximately at $Eo > 200$. Numerical solutions with either wavy or breaking up bubble wakes and cases with small bubble detachment required extremely refined meshes to be accurate. Fig. 18 illustrates an example for this unstable nature of gas-liquid interface for case 36 in Table 1.

The wake region is investigated by two main parameters, namely dimensionless wake length (L_w/D), and dimensionless wake volume (V_w/D^3). One of the main correlations used to predict these parameter is that of Campos and De Carvalho (1988) which is given by equation (15). Recently, Araújo et al. (2012) suggested equation for the dimensionless wake length, L_w/D , and dimensionless wake volume (V_w/D^3) function in Eo , and M , given by:

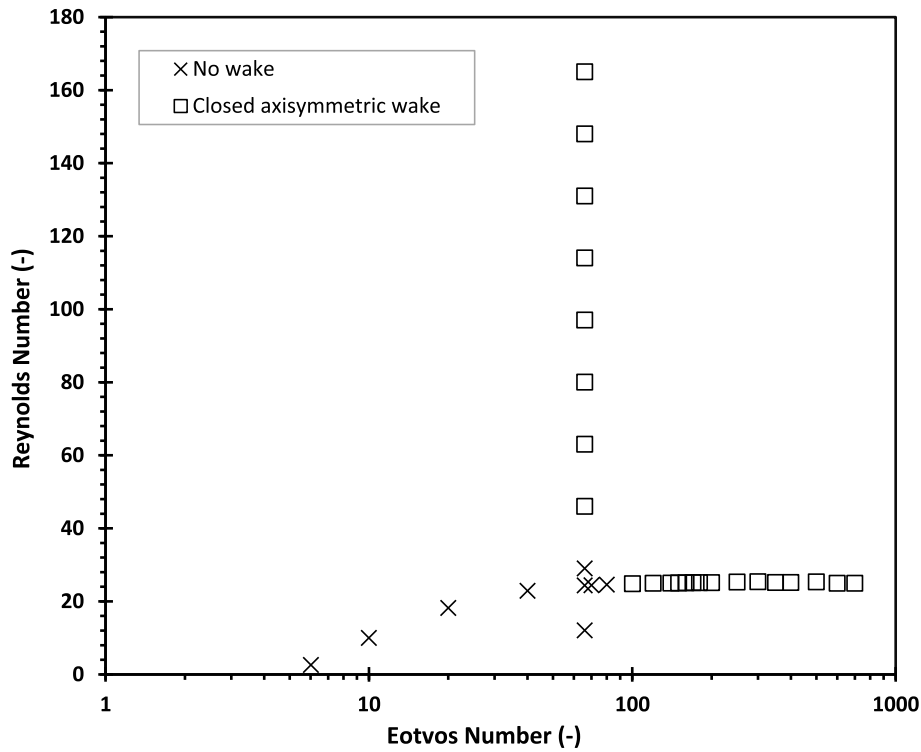


Fig. 16. Phase diagram of the presence and kind of wake of Taylor bubbles rising through vertical columns of stagnant liquid.

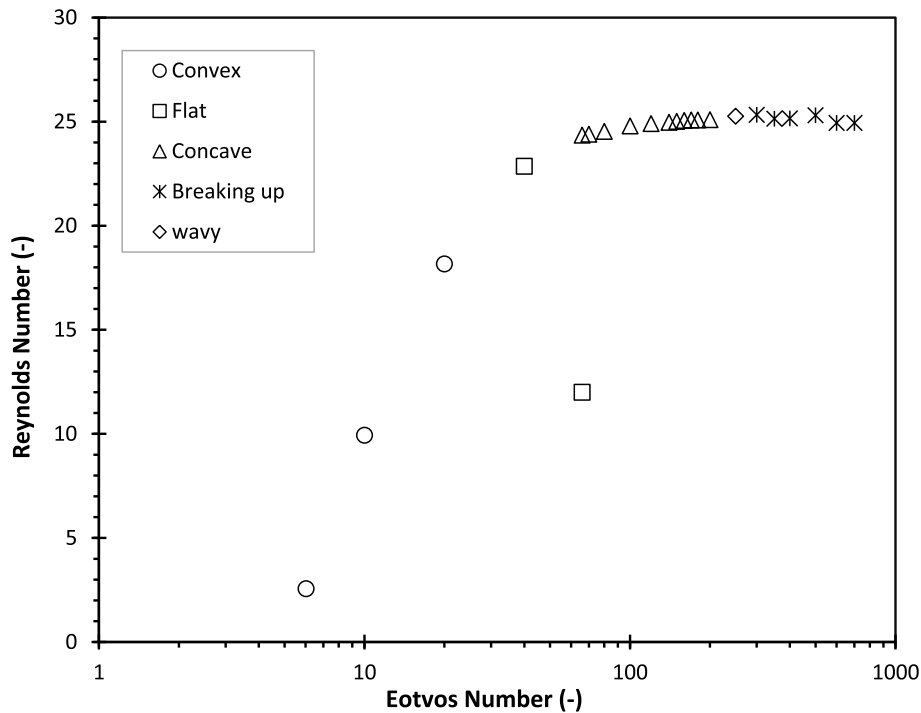


Fig. 17. Phase diagram of the rear of Taylor bubble shape, expressed in terms of Eo , and $Re_{U_{TB}}$, of a Taylor bubbles rising through vertical columns of stagnant liquid.

$$\frac{L_W}{D} = \frac{1}{4} \times [0.555 - 7.793 \times 10^3 \ln(M)] \times \ln\left(\frac{Eo^3}{M}\right) - 2.133 + 8.046 \times 10^{-2} \ln(M)$$

$$\frac{V_W}{D^3} = 1.365 \times 10^{-1} \left(\frac{L_W}{D}\right)^2 + 2.176 \times 10^{-1} \left(\frac{L_W}{D}\right) - 2.91910^{-2}$$

(20)

Following the same procedure of Araujo et al. (2012) for measuring

L_W/D , and V_W/D^3 , Table 1 gives the simulation results for these two parameters for all cases under investigation.

As discussed earlier, the effect of $Re_{U_{TB}}$ on the flow in the wake region is noticed by the change in the shape of rear of the Taylor bubble from flat into convex. In addition, the increase in $Re_{U_{TB}}$ decreases the liquid film thickness thus squeezing liquid in narrower region. As a result, the intensity and size of the wake increases, which is noticed by the increment

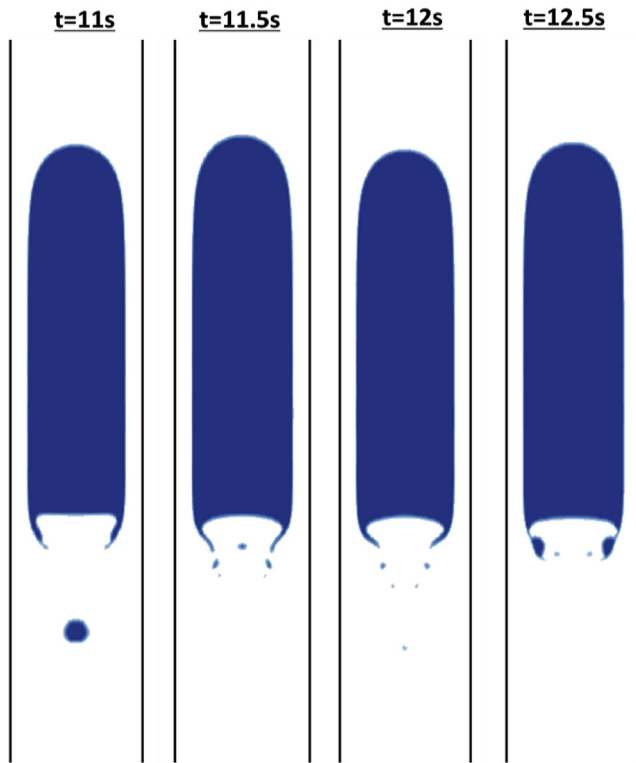


Fig. 18. Development of wake flow pattern of Taylor bubble rising vertical columns of stagnant liquid for case (36) with a time interval of 0.5s: $Eo = 250$, $Re_{U_{TB}} = 25.26$, and $Fr_{U_{TB}} = 0.3$.

of L_W/D and V_W/D^3 as indicated in Fig. 19. Furthermore, A similar scenario is noticed for the increase in Eo , for cases 21 to 42 in Table 1, that results in increase in L_W/D , and V_W/D^3 values. It should be pointed out that for cases 36 to 42 in Table 1, L_W/D and V_W/D^3 are calculated as average values once the solution is converged and most of the hydrodynamics characteristics investigated are developed. This is because the

developed bubble shape is unstable which is either wavy or breaking up. A good matching between simulation results for L_W/D , and V_W/D^3 , and the experimental correlation given by Campos and De Carvalho (1988) is shown in Fig. 19.

Also, Table 4 gives a comparison of the simulation results, experimental data, and correlations from literature for L_W/D , and V_W/D^3 for validation of the present code regarding flow in wake region. Two simulation cases based on the experimental work of Campos and De Carvalho (1988) with Eo in range of 64, and $Re_{U_{TB}}$ in range of 56–107, are simulated. It can be concluded from Table 4, that numerical results show good matching with other correlations in literature, experimental correlation of Campos and De Carvalho (1988) and numerical correlation of Araújo et al. (2012).

In conclusion, the flow in the wake region is significantly affected by both Eo , and $Re_{U_{TB}}$ which can be clarified by the change in corresponding

Table 4

Numerical and experimental values of L_W/D , and V_W/D^3 , and their respective deviations.

1. Experimental data ^a				
	$\frac{L_W}{D}(-)$	Error (%)	$\frac{V_W}{D^3}(-)$	Error (%)
Case 1	0.5145	–	0.1319	–
Case 2	0.544	–	0.147	–
2. Current simulation				
	$\frac{L_W}{D}(-)$	Error (%)	$\frac{V_W}{D^3}(-)$	Error (%)
Case 1	0.487	5.34	0.1094	17.06
Case 2	0.5567	-2.33	0.1394	5.17
3. Experimental correlation ^a				
	$\frac{L_W}{D}(-)$	Error (%)	$\frac{V_W}{D^3}(-)$	Error (%)
Case 1	0.5145	–	0.0973	26.23
Case 2	0.5498	-1.07	0.1227	16.53
4. Numerical correlation ^b				
	$\frac{L_W}{D}(-)$	Error (%)	$\frac{V_W}{D^3}(-)$	Error (%)
Case 1	0.40616	21.06	0.0817	38.06
Case 2	0.4754	12.61	0.1051	28.50

^a Campos and De Carvalho (1988).

^b Araújo et al. (2012).

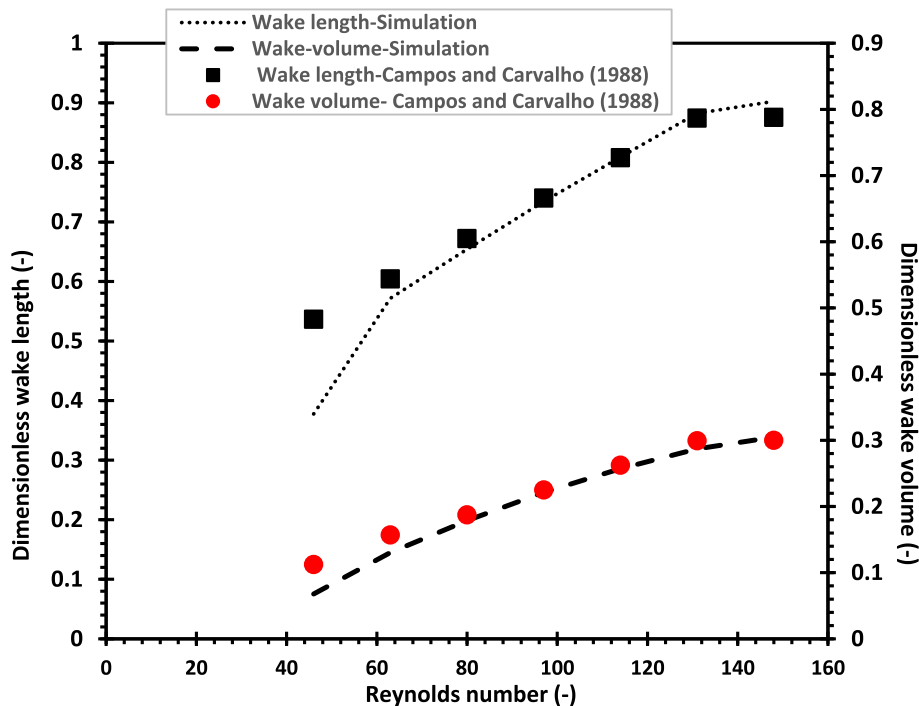


Fig. 19. Numerical results of dimensionless wake length (L_W/D), and dimensionless wake volume (V_W/D^3) for cases 13 to 19 in Table 1.

dimensionless wake length and volume. Additionally, the results revealed that both L_W/D and V_W/D^3 are mainly dependent on Eo , and $Re_{U_{TB}}$, which again agrees with the predictions of Kang et al. (2010), and contradicts Araújo et al. (2012) conclusion with different point of view based on the significance of Reynolds number ($Re_{U_{TB}}$) rather than Morton number (M).

5. Conclusions

In this paper, a complete dimensionless analysis of single Taylor bubble rising through a vertical stagnant Newtonian liquid problem is performed followed by the order of magnitude analysis of the equations of motion. The main conclusion is that Froude number ($Fr_{U_{TB}}$), is function of Eötvös number (Eo), Reynolds number ($Re_{U_{TB}}$), density ratio (Γ_ρ), viscosity ratio (Γ_μ) and L_{TB}/D . The effect of density ratio (Γ_ρ), viscosity ratio (Γ_μ) and the bubble size to the pipe diameter (L_{TB}/D) are examined for the sake of supporting other previous experimental and numerical works in the literature.

Based on the dimensionless analysis, the hydrodynamic characteristics of single Taylor bubble rising through vertical stagnant Newtonian liquid is investigated using the volume-of-fluid (VOF) methodology implemented in the computational fluid dynamics software package, ANSYS Fluent (Release 15.0), with particular focus on the sole dimensionless parameters: Froude number ($Fr_{U_{TB}}$), Reynolds number ($Re_{U_{TB}}$) and Eötvös number (Eo). The results are divided into four sections: Taylor bubble shape, Taylor bubble rise velocity, flow in the liquid film region and flow in the wake region. The numerical results predicted the problem effectively as they were validated by testing some of the selective cases against both theoretical and experimental data in the literature.

Based on the numerical results and using guidelines from order of magnitude analysis, correlation to estimate Taylor bubble rise velocity (U_{TB}) as a function in only Eo , and $Re_{U_{TB}}$ is proposed and shows good prediction when compared with other well-known correlations in literature, especially with Viana et al. (2003) correlation.

Detailed hydrodynamics

Eo showed significant effect on all hydrodynamic features of Taylor bubbles. The increment in surface tension increases the curvature of the bubble nose, increases the liquid film thickness, decreases the dimen-

sionless velocity of fully developed falling liquid film, and finally increases wake intensity and size. For the flow in the liquid film, the numerical results are compared with the theoretical predictions given by Brown (1965) and establishing a strong correspondence.

$Re_{U_{TB}}$ shows contribution similar to Eo number. The developed Taylor bubble shape changes from long slender shape into shorter and flatter bubbles with the increase in $Re_{U_{TB}}$; this results in reduction in liquid film thickness.

Additionally, the wall shear stress was examined under the effect of both Eo and $Re_{U_{TB}}$. The increase in $Re_{U_{TB}}$ results in less viscous effect in the liquid surrounding the Taylor bubble that subscribes to decrement in wall shear stress. The wall shear stress increases with Eo , however it should be pointed out that lower values of Eo possesses different distribution for the wall shear stress with noticeable peak value right before the Taylor bubble rear due to the sharp flat and convex shape of the bubble attained at lower values of Eo . Based on the numerical results for wall shear stress, a proposed correlation for maximum wall shear stress is developed and predicts results favourably matching with the theoretical predictions of Brown (1965).

Furthermore, for the wake region, a phase diagram showing the presence and nature of wake, and of the shape of the rear Taylor bubble is illustrated. Unstable bubbles are developed at Eo above 200 which are characterized by presence of small bubbles shedding off from main Taylor bubble into wake region.

Finally, it can be concluded that the developed numerical results agree well with the order of magnitude analysis. For instance, the order of magnitude analysis has shown that, in order for Reynolds number for the major viscous terms to remain intact, it should be of order $\left(\frac{L_{TB}}{D}\right)$, which means relatively large values. The numerical result agrees well with that showing the significance of the larger values of Reynolds number on the dynamics of Taylor bubble including bubble shape, terminal velocity, flow in the liquid film and flow in the wake region.

Acknowledgements

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Appendix 1. Order of magnitude analysis for the equation of motion

The Navier-stokes equations in polar coordinates assuming axisymmetric flow are:

$$\rho_L \left(\frac{\partial v_r}{\partial t} + v_r \frac{\partial v_r}{\partial r} + v_z \frac{\partial v_r}{\partial z} \right) = \rho_L g_r - \frac{\partial p}{\partial r} + \mu \left[\frac{\partial^2 v_r}{\partial r^2} + \frac{1}{r} \frac{\partial v_r}{\partial r} + \frac{\partial^2 v_r}{\partial z^2} - \frac{v_r}{r^2} \right] \tag{1-1}$$

$$\rho_L \left(\frac{\partial v_z}{\partial t} + v_r \frac{\partial v_z}{\partial r} + v_z \frac{\partial v_z}{\partial z} \right) = \rho_L g_z - \frac{\partial p}{\partial z} + \mu \left[4 \frac{\partial^2 v_z}{\partial r^2} + \frac{1}{r} \frac{\partial v_z}{\partial r} + \frac{\partial^2 v_z}{\partial z^2} \right] \tag{1-2}$$

$$0 = -\frac{1}{r} \frac{\partial p}{\partial \theta} \quad (v_\theta = 0) \tag{1-3}$$

And the continuity equation is:

$$\frac{\partial v_r}{\partial r} + \frac{v_r}{r} + \frac{\partial v_z}{\partial z} = 0 \tag{1-4}$$

Introducing the following dimensionless variables:

$$v_r^* = \frac{v_r}{U_{TB}}, \quad v_z^* = \frac{v_z}{U_{TB}}, \quad t^* = t \left(\frac{U_{TB}}{L_{TB}} \right), \quad r^* = \frac{r}{(D/2)}, \quad z^* = \frac{z}{L_{TB}}, \quad p^* = \frac{p}{\frac{1}{2} \rho_L U_{TB}^2}, \quad (1-5)$$

$$g_r^* = \frac{g_r}{g}, \quad g_z^* = \frac{g_z}{g}, \quad \sigma_r^* = \frac{\sigma_r}{\sigma}, \quad \sigma_z^* = \frac{\sigma_z}{\sigma}, \quad \sigma_\theta^* = \frac{\sigma_\theta}{\sigma}, \quad k^* = \frac{k}{(1/D^2)}$$

The Navier-stokes equations and continuity equations in dimensionless form can be summarized as follows:

$$\left(\frac{D}{L_{TB}} \right) \frac{\partial v_r^*}{\partial t^*} + 2 v_r^* \frac{\partial v_r^*}{\partial r^*} + \left(\frac{D}{L_{TB}} \right) v_z^* \frac{\partial v_r^*}{\partial z^*} = \frac{1}{Fr_{U_{TB}}^2} g_r^* - \frac{\partial p^*}{\partial r^*} + \frac{1}{Re_{U_{TB}}} \left[4 \frac{\partial^2 v_r^*}{\partial r^{*2}} + \frac{4}{r^*} \frac{\partial v_r^*}{\partial r^*} + \left(\frac{D}{L_{TB}} \right)^2 \frac{\partial^2 v_r^*}{\partial z^{*2}} - \frac{4v_r^*}{r^*} \right] \quad (1-6)$$

$$\left(\frac{D}{L_{TB}} \right) \frac{\partial v_z^*}{\partial t^*} + 2 v_z^* \frac{\partial v_z^*}{\partial r^*} + \left(\frac{D}{L_{TB}} \right) v_z^* \frac{\partial v_z^*}{\partial z^*} = \frac{1}{Fr_{U_{TB}}^2} g_z^* - \left(\frac{D}{L_{TB}} \right) \frac{\partial p^*}{\partial z^*} + \frac{1}{Re_{U_{TB}}} \left[4 \frac{\partial^2 v_z^*}{\partial r^{*2}} + \frac{4}{r^*} \frac{\partial v_z^*}{\partial r^*} + \left(\frac{D}{L_{TB}} \right)^2 \frac{\partial^2 v_z^*}{\partial z^{*2}} \right] \quad (1-7)$$

$$0 = \frac{\partial p^*}{\partial \theta^*} \quad (1-8)$$

$$2 \left(\frac{\partial v_r^*}{\partial r^*} + \frac{v_r^*}{r^*} \right) + \left(\frac{D}{L_{TB}} \right) \frac{\partial v_z^*}{\partial z^*} = 0 \quad (1-9)$$

where $Fr_{U_{TB}} = \frac{U_{TB}}{\sqrt{gD}}$, $Re_{U_{TB}} = \frac{\rho_L U_{TB} D}{\mu_L}$.

On the gas-liquid interface, the following conditions are applied:

$$\left[-\frac{\partial p^*}{\partial r^*} + \frac{1}{Fr_{U_{TB}}^2} k^* \sigma_r^* \right]_S + \frac{4}{Re_{U_{TB}}} \left[\frac{\partial^2 v_r^*}{\partial r^{*2}} + \frac{1}{r^*} \frac{\partial v_r^*}{\partial r^*} \right]_S = 0 \quad (1-10)$$

$$\left[\left(\frac{D}{L_{TB}} \right) \frac{\partial v_z^*}{\partial t^*} + \left(\frac{D}{L_{TB}} \right) \frac{\partial v_z^*}{\partial r^*} + \left(\frac{D}{L_{TB}} \right) v_z^* \frac{\partial v_z^*}{\partial z^*} \right]_S = \left[-\left(\frac{D}{L_{TB}} \right) \frac{\partial p^*}{\partial z^*} + \frac{1}{Re_{U_{TB}}} \left[4 \frac{\partial^2 v_z^*}{\partial r^{*2}} + \frac{4}{r^*} \frac{\partial v_z^*}{\partial r^*} + \left(\frac{D}{L_{TB}} \right)^2 \frac{\partial^2 v_z^*}{\partial z^{*2}} \right] \right]_S + \frac{1}{Fr_{U_{TB}}^2} \left[g_z^* + \frac{1}{Eo} k^* \sigma_z^* \right]_S \quad (1-11)$$

$$0 = \left[-\frac{1}{r^*} \frac{\partial p^*}{\partial \theta^*} \right]_S + \frac{1}{Fr_{U_{TB}}^2} \left[k^* \sigma_\theta^* \right]_S \quad (1-12)$$

where $Eo = \frac{g \rho_L D^2}{\sigma}$.

Performing similar analysis with respect to gas and liquid sides would show the additional dimensionless groups: density ratio (Γ_ρ), viscosity ratio (Γ_μ).

In order to perform order of magnitude analysis to the continuity equation, the following orders are introduced to equations (1-9):

$$r^* = o(1), z^* = o(1), v_z^* = o(1), k^* = o(1), \sigma_z^* = o(1), g_z^* = o\left(\frac{D}{L_{TB}}\right) \quad (1-13)$$

$$g_r^* = o\left[\frac{D}{L_{TB}}\right]^2 \text{ and } t^* = o(1)$$

It can be concluded from equations (1-9) and (1-13), in order to keep the continuity equation intact v_r^* should be of order $\frac{D}{L_{TB}}$, that is $v_r^* = o\left(\frac{D}{L_{TB}}\right)$.

Applying the same analysis to momentum equation in the z^* direction (equations (1-7)), will give the following:

$$Fr_{U_{TB}} = o(1), \quad Re_{U_{TB}} = o\left(\frac{L_{TB}}{D}\right) \quad (1-14)$$

Applying the same analysis to the momentum equation in r^* direction (equations (1-6)) and using the conclusion given in equations (1-14), will give the following:

$$\frac{\partial p^*}{\partial r^*} = o\left(\frac{D}{L_{TB}}\right)^2, \quad p^* = o\left(\frac{D}{L_{TB}}\right)^2, \quad \therefore \frac{\partial p^*}{\partial z^*} = o\left(\frac{D}{L_{TB}}\right)^2 \text{ - which is very small} \quad (1-15)$$

Finally, applying the same analysis to the gas-liquid interface (equations (1-11)), gives the following:

$$Eo = o\left(\frac{L_{TB}}{D}\right) \quad (1-16)$$

It is worth noting that, if the characteristic velocity U_{TB} is placed by a new characteristic velocity $V_s = (gD)^{1/2}$, the dimensionless variables given in

equations (1–5) will be modified as follows;

$$\hat{v}_r = \frac{v_r}{V_s}, \quad \hat{v}_z = \frac{v_z}{V_s}, \quad \hat{t} = t \left(\frac{V_s}{L_{TB}} \right), \quad \hat{p} = \frac{p}{\frac{1}{2} \rho_L V_s^2} \quad (1-17)$$

Thus for instance, the dimensionless momentum equation in the radial direction (equations (1–6)) will be as follows:

$$\left(\frac{D}{L_{TB}} \right) \frac{\partial \hat{v}_r}{\partial \hat{t}} + 2 \hat{v}_r \frac{\partial \hat{v}_r}{\partial r^*} + \left(\frac{D}{L_{TB}} \right) \hat{v}_z \frac{\partial \hat{v}_r}{\partial z^*} = g_r^* - \frac{\partial \hat{p}}{\partial r^*} + \frac{1}{Re_{modified}} \left[4 \frac{\partial^2 \hat{v}_r}{\partial r^{*2}} + \frac{4}{r^*} \frac{\partial \hat{v}_r}{\partial r^*} + \left(\frac{D}{L_{TB}} \right)^2 \frac{\partial^2 \hat{v}_r}{\partial z^{*2}} - \frac{4 \hat{v}_r}{r^{*2}} \right] \quad (1-18)$$

where $Re_{modified}$ is the modified Reynolds number which is defined as follows:

$$Re_{modified} = \frac{V_s D}{\nu} = \frac{(gD)^{1/2} D}{\nu} = \frac{(gD^3)^{1/2}}{\nu} = N_f \quad (1-19)$$

It can be concluded from equations (1–18) that Froude number disappeared from the dimensionless governing equation as being of order unity as shown in equations (1–14). Hence, in this frame with the new definition of the characteristic velocity (V_s), the inverse viscosity number (N_f) can represent the ratio of inertia force to viscous force and Froude number is no longer a dimensionless parameter.

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